DYNAMICS IN DUMBBELL DOMAINS III. CONTINUITY OF ATTRACTORS

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ABSTRACT. In this paper we conclude the analysis started in [3] and continued in [4] concerning the behavior of the asymptotic dynamics of a dissipative reactions diffusion equation in a dumbbell domain as the channel shrinks to a line segment. In [3], we have established an appropriate functional analytic framework to address this problem and we have shown the continuity of the set of equilibria. In [4], we have analyzed the behavior of the limiting problem. In this paper we show that the attractors are upper semicontinuous and, moreover, if all equilibria of the limiting problem are hyperbolic, then they are lower semicontinuous and therefore, continuous. The continuity is obtained in L^p and H^1 norms.

1. Introduction

This paper is concerned with the continuity of the asymptotic dynamics of a dissipative reaction-diffusion equation in a dumbbell type domain as the channel degenerates to a line segment. Here we conclude the analysis started in [3], where we studied the continuity of the equilibria, and continued in [4], where we studied the limiting problem. We refer to the introduction in [3] for a broad perspective of the problem.

More precisely, we consider a reaction-diffusion equation of the form

$$\begin{cases} u_t - \Delta u + u = f(u) & x \in \Omega_{\varepsilon} \\ \frac{\partial u}{\partial n} = 0, & x \in \partial \Omega_{\varepsilon} \end{cases}$$
 (1.1)

where, for $N \geq 2$ and $\varepsilon \in (0,1]$, $\Omega_{\varepsilon} \subset \mathbb{R}^N$ is a typical dumbbell domain; that is, two disconnected domains, denoted by Ω , joined by a thin channel, denoted by R_{ε} . The channel R_{ε} degenerates to a line segment as the parameter ε approaches zero, see Figure 1. We refer to [3] Section 2, for a complete and rigorous definition of the dumbbell domain that we are

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considering. We mention that the channels R_{ε} considered here are fairly general and are not required to be cylindrical.

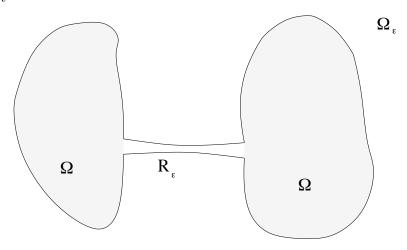


Figure 1. Dumbbell domain

The limit "domain" consists of the fixed part Ω and the line segment R_0 . Without loss of generality, we may assume that $R_0 = \{(x, 0, \dots, 0) : 0 < x < 1\}$, see Figure 2 of [4].

The limit equation is given by

$$\begin{cases} w_t - \Delta w + w = f(w), & x \in \Omega, \ t > 0 \\ \frac{\partial w}{\partial n} = 0, & x \in \partial \Omega \end{cases}$$

$$\begin{cases} v_t - \frac{1}{g} (gv_x)_x + v = f(v), & x \in (0, 1) \\ v(0) = w(P_0), \ v(1) = w(P_1) \end{cases}$$

$$(1.2)$$

where w is defined in Ω , v is defined in R_0 and P_0 , P_1 are the points where the line segment touches the boundary of Ω . Observe that the boundary conditions of v in (0,1) are given in terms of a continuity condition, so that the whole function (w,v) is continuous in the junction between Ω and R_0 . The function $g:[0,1] \to (0,\infty)$ is a smooth function related to the geometry of the channel R_{ε} , more exactly, on the way the channel R_{ε} collapses to the segment line R_0 , see [3]. For instance, if the channel is given by $R_{\varepsilon} = \{(x, \varepsilon x') : (x, x') \in R_1\}$, for some fixed reference channel R_1 , then $g(x) = |\{x' : (x, x') \in R_1\}|_{N-1}$, where $|\cdot|_{N-1}$ is the (N-1)-dimensional Lebesgue measure, see [3].

In [3] we have studied how the equilibria of (1.1) behave as the parameter ε tends to zero. Since the spaces to which the equilibria belong also vary with ε , we developed an appropriate functional analytical setting to compare these functions as well as deal with this singular perturbation problem. We have constructed the family of spaces U_{ε}^{p} , $0 < \varepsilon \leq 1$, in Ω_{ε} , which is the space $L^{p}(\Omega_{\varepsilon})$ with the norm

$$||u_{\varepsilon}||_{U_{\varepsilon}^p}^p = \int_{\Omega} |u|^p + \frac{1}{\varepsilon^{N-1}} \int_{R_{\varepsilon}} |u_{\varepsilon}|^p.$$

Observe that the integral in R_{ε} has the weight $1/\varepsilon^{N-1}$, which amplifies the effect of a function in the channel. As observed in [3] a constant function in R_{ε} will converge to zero if we do not introduce the appropriate weight $(1/\varepsilon^{N-1})$. In this setting, we showed that the appropriate limit space should be $U_0^p = L^p(\Omega) \oplus L_g^p(0,1)$; that is, $(w,v) \in U_0^p$ iff $w \in L^p(\Omega)$, $v \in L^p(0,1)$. The norm in U_0^p is given by

$$\|(w,v)\|_{U_0^p}^p = \int_{\Omega} |w|^p + \int_0^1 g|v|^p.$$

If $A_{\varepsilon}: D(A_{\varepsilon}) \subset U_{\varepsilon}^p \to U_{\varepsilon}^p$ is given by $A_{\varepsilon}(u) = -\Delta u + u$ for $0 < \varepsilon \leqslant 1$, and $A_0: D(A_0) \subset U_0^p \to U_0^p$ is given by $A_0(w,v) = (-\Delta u + u, -\frac{1}{g}(gv_x)_x + v)$, we proved in Proposition 2.7 of [3] that $A_{\varepsilon}^{-1} \xrightarrow{\varepsilon \to 0} A_0^{-1}$. Moreover, considering the equilibria of (1.1) and (1.2), in an abstract way, as the solutions of

$$A_{\varepsilon}u = F_{\varepsilon}(u), \ \varepsilon \in [0,1],$$

with F_{ε} being suitable Nemitskiĭ maps, or as fixed points of the nonlinear maps $A_{\varepsilon}^{-1} \circ F_{\varepsilon}$: $U_{\varepsilon}^{p} \to U_{\varepsilon}^{p}$, we showed the convergence of the equilibria see Theorem 2.3 of [3]. Also, if the equilibria of the limiting problem (1.2) are hyperbolic, we proved the convergence of the resolvent of linearizations around the equilibria and the convergence of the linear unstable manifolds.

In [4] we studied in detail the properties of the limiting problem in terms of generation of linear singular semigroups by the operator A_0 , local well posedness and existence of attractor for the associated singular nonlinear semigroup. We also show that, when all equilibria are hyperbolic, the attractor of the limiting problem (which is not gradient) can be characterized as the union of the unstable manifolds of the equilibria.

As we mentioned in the introduction of [3], our final objective is to compare the whole dynamics of problems (1.1) and (1.2). That is, to prove the continuity of the attractors as ε tends to zero. To accomplish this goal, we proposed an agenda based on a deep and thorough study of the linear part of the problems consisting on the study of the convergence properties of the resolvent operators. That agenda was established in the introduction of [3] and consisted of six items. The first three were covered in [3].

In this paper we consider the last three items of that agenda and complete the analysis. Hence, we show the convergence of the resolvent operators $(\lambda + A_{\varepsilon})^{-1}$ to $(\lambda + A_0)^{-1}$ and use this information to obtain the convergence of the linear semigroups. With the variation of constants formula and the convergence of linear semigroups we show the convergence of the nonlinear semigroups, from which the upper semicontinuity of the attractors follows easily. This is done in a very similar manner as in [2].

Finally, if each equilibria of the limitting problem is hyperbolic, with the convergence of the equilibria and of its linear unstable manifolds, we show the convergence of the local nonlinear unstable manifolds of equilibria. Using the gradient-like structure of the limiting equation we prove lower semicontinuity (and therefore the continuity) of the attractors.

Next, we describe contents of the paper. In Section 2 we recall the general setting of the problem and state the main results of this paper; that is, the upper and lower semicontinuity of the attractors. In Section 3 we study the convergence of the resolvent operators associated with the linear operators obtaining rates of convergence of equilibria and of resolvent

operators associated to the linearizations around equilibria. Based in the resolvent estimates obtained in Section 3, we analyze in Section 4 the convergence of the linear semigroups. In Section 5 we obtain the continuity of the nonlinear semigroups and the upper semicontinuity of the attractors. In Section 6 and, under the assumption that all equilibria of the limiting problem are hyperbolic, prove that the local unstable manifolds behave continuously as ε tends to zero. The continuity of local unstable manifolds is the key step to show the continuity of the attractors. Finally in Section 7, we analyze the continuity properties of the attractors in other norms.

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Special dedication. The question of the continuity of attractors of reaction-diffusion equations in dumbbell domains, as it is addressed in this paper as well as in [3, 4], was raised by Jack K. Hale and a great amount of the ideas and techniques explored in the three articles were proposed initially by him. The authors are specially grateful for his permanent support and motivation and would like to dedicate this work to him on the occasion of his 80th birthday.

2. Setting of the problem and statement of the main results

The setting is the same as the one we established initially in [3]. We recall some of the terminology which will be needed to study the continuity of attractors.

Consider the spaces U^p_{ε} and U^p_0 defined in Section 1, see also [3]. Let $0 < \varepsilon \leqslant 1$ and let $A_{\varepsilon}: D(A_{\varepsilon}) \subset U^p_{\varepsilon} \to U^p_{\varepsilon}$, $1 \leqslant p < \infty$, be the linear operator defined by

$$D(A_{\varepsilon}) = \left\{ u \in W^{2,p}(\Omega_{\varepsilon}) : \Delta u \in U_{\varepsilon}^{p}, \, \partial u / \partial n = 0 \text{ in } \partial \Omega_{\varepsilon} \right\},$$

$$A_{\varepsilon}u = -\Delta u + u, \ u \in D(A_{\varepsilon}). \tag{2.1}$$

Also, for $p > \frac{N}{2}$, let $A_0: D(A_0) \subset U_0^p \to U_0^p$ be the operator defined by

$$D(A_0) = \{(w, v) \in U_0^p : w \in D(\Delta_N^{\Omega}), (gv')' \in L^p(0, 1), v(0) = w(P_0), v(1) = w(P_1)\}$$
 (2.2)

$$A_0(w,v) = \left(-\Delta w + w, -\frac{1}{q}(gv')' + v\right), \quad (w,v) \in D(A_0),$$
(2.3)

where Δ_N^{Ω} is the Laplace operator with homogeneous Neumann boundary conditions in $L^p(\Omega)$ with $D(\Delta_N^{\Omega}) = \{u \in W^{2,p}(\Omega) : \frac{\partial u}{\partial n} = 0 \text{ in } \partial\Omega\}.$

We note that, for $p > \frac{N}{2}$ we have that $D(\Delta_N^{\Omega})$ is continuously embedded in $C(\bar{\Omega})$. In that case, the functions in $D(\Delta_N^{\Omega})$ have well defined traces at P_0 and P_1 .

Recall that we have defined in [3] the operator $M_{\varepsilon}: U_{\varepsilon}^p \to U_0^p$, as follows

$$\psi_{\varepsilon} \to (M_{\varepsilon}\psi_{\varepsilon})(z) = \begin{cases} \psi_{\varepsilon}(z), & z \in \Omega \\ \frac{1}{|\Gamma_{\varepsilon}^{z}|} \int_{\Gamma_{\varepsilon}^{z}} \psi(z, y) dy, & z \in (0, 1), \end{cases}$$
 (2.4)

where $\Gamma_{\varepsilon}^z = \{y : (z,y) \in R_{\varepsilon}\}$. It is easy to see, from Fubini-Tonelli Theorem and Hölder inequality, that M_{ε} is a well defined bounded linear operator with $\|M_{\varepsilon}\|_{\mathcal{L}(U_{\varepsilon}^p,U_0^p)} = 1$.

Also consider the family of extension operators $E_{\varepsilon}: U_0^p \to U_{\varepsilon}^p$ defined by

$$E_{\varepsilon}(w,v)(x) = \begin{cases} w(x), & x \in \Omega \\ v(s), & (s,y) \in R_{\varepsilon}. \end{cases}$$
 (2.5)

It is very easy to see that $||E_{\varepsilon}(w,v)||_{U_{\varepsilon}^{p}} = ||(w,v)||_{U_{0}^{p}}$.

The operator A_{ε} generates an analytic semigroup $\{e^{A_{\varepsilon}t}: t \geq 0\}$ on U_{ε}^p whereas, from the results in [4], the operator A_0 generates a *singular semigroup* in U_0^p that we will denote by $\{e^{-A_0t}: t \geq 0\}$, see [4].

We rewrite (1.1) and (1.2) in the abstract form

$$\begin{cases}
\dot{u}_{\varepsilon} + A_{\varepsilon} u_{\varepsilon} = f_{\varepsilon}(u_{\varepsilon}) \\
u_{\varepsilon}(0) = u_{0}^{\varepsilon} \in U_{\varepsilon}^{p}
\end{cases}$$
(2.6)

and

$$\begin{cases} \dot{u} + A_0 u = f_0(u) \\ u(0) = u_0 \in U_0^p \end{cases}$$
 (2.7)

With respect to the nonlinearity f, we will assume that

- (i) $f: \mathbb{R} \to \mathbb{R}$ is a C^2 function,
- (ii) $|f(u)| + |f'(u)| + |f''(u)| \leqslant C_1$ for all $u \in \mathbb{R}$.

Remark 2.1. From the point of view of studying the asymptotic dynamics (continuity of attractors), the assumption (ii) does not imply any restriction on the nonlinearities. Since we are assuming that f is dissipative, under the usual growth assumptions, the attractors are bounded in $L^{\infty}(\Omega_{\varepsilon})$ uniformly with respect to $\varepsilon \in [0,1]$ and one may cut the nonlinearities to make them satisfy the above assumptions (See Remark 2.2 of [3]).

Under these assumptions, the nonlinear semigroups $\{T_{\varepsilon}(t): t \geq 0\}$ in U^p_{ε} associated with (2.6) and the singular semigroup $\{T_0(t): t \geq 0\}$ in U^p_0 , p > N/2, associated with (2.7), have compact global attractors $\mathcal{A}_{\varepsilon} \subset U^p_{\varepsilon}$ and $\mathcal{A}_0 \subset U^p_0$ respectively (see [4]). In general, the attractors lie in more regular spaces and in particular, from comparison arguments, they lie in U^{∞}_{ε} and U^{∞}_{0} .

The following concept of E-convergence has been proved to be very appropriate when dealing with sequences of functions in different spaces, see [14, 7, 3].

Definition 2.2. We say that a sequence $\{u_{\varepsilon}\}_{{\varepsilon}\in(0,1]}$, $u_{\varepsilon}\in U_{\varepsilon}^p$, E_{ε} -converges to $u_0\in U_0^p$ if $\|u_{\varepsilon}-E_{\varepsilon}u_0\|_{U_{\varepsilon}^p}\stackrel{{\varepsilon}\to 0}{\longrightarrow} 0$ (see (2.5) for the definition of E_{ε}). We write this as $u_{\varepsilon}\stackrel{E}{\longrightarrow} u_0$.

This notion of convergence can be extended to sets in the following manner (see [7]).

Definition 2.3. Let $\mathcal{A}_{\varepsilon} \subset U_{\varepsilon}^p$, $\varepsilon \in [0,1]$ and $\mathcal{A}_0 = \mathcal{A} \subset U_0^p$. Denote by $\operatorname{dist}(\cdot, \cdot)$ the metric induced by the norm in U_{ε}^p , $\varepsilon \in [0,1]$, i.e. $\operatorname{dist}(u_{\varepsilon}, v_{\varepsilon}) = \|u_{\varepsilon} - v_{\varepsilon}\|_{U_{\varepsilon}^p}$.

- (1) We say that the family of sets $\{A_{\varepsilon}\}_{{\varepsilon}\in[0,1]}$ is E_{ε} -upper semicontinuous at ${\varepsilon}=0$ if $\sup_{u_{\varepsilon}\in A_{\varepsilon}} \operatorname{dist}(u_{\varepsilon}, E_{\varepsilon}A) \stackrel{{\varepsilon}\to 0}{\longrightarrow} 0$. (2) We say that the family of sets $\{A_{\varepsilon}\}_{{\varepsilon}\in[0,1]}$ is E_{ε} -lower semicontinuous at ${\varepsilon}=0$ if
- (2) We say that the family of sets $\{A_{\varepsilon}\}_{{\varepsilon}\in[0,1]}$ is E_{ε} -lower semicontinuous at ${\varepsilon}=0$ if $\sup_{u\in\mathcal{A}}\operatorname{dist}(E_{\varepsilon}u,\mathcal{A}_{\varepsilon})\stackrel{{\varepsilon}\to 0}{\longrightarrow} 0.$

Remark 2.4. In order to show the upper or lower semicontinuity of sets, the following characterizations are usefull

- (1) If any sequence $\{u_{\varepsilon}\}$ with $u_{\varepsilon} \in \mathcal{A}_{\varepsilon}$ has an E_{ε} -convergent subsequence with limit belonging to \mathcal{A} , then $\{\mathcal{A}_{\varepsilon}\}$ is E_{ε} -upper semicontinuous at zero.
- (2) If \mathcal{A} is compact and for any $u \in \mathcal{A}$ there is a sequence $\{u_{\varepsilon}\}$ with $u_{\varepsilon} \in \mathcal{A}_{\varepsilon}$, which E_{ε} -converges to u, then $\{\mathcal{A}_{\varepsilon}\}$ is E_{ε} -lower semicontinuous at zero.

With all this concepts in mind, our main result is the following,

Theorem 2.5. The family of attractors $\{A_{\varepsilon}\}_{{\varepsilon}\in[0,1]}$ is E_{ε} -upper semicontinuous at ${\varepsilon}=0$ in U_{ε}^p for every $1\leqslant p<\infty$.

Moreover, if every equilibria of the limit problem is hyperbolic, then the family of attractors is also E_{ε} - lower semicontinuous at $\varepsilon = 0$ in U_{ε}^p for every $1 \leq p < \infty$.

Remark 2.6. Observe that once the statement of Theorem 2.5 is shown for a particular $p \geqslant 1$, then from the boundedness of the attractors in U_{ε}^{∞} and U_{0}^{∞} , it will also be proved for all $1 \leqslant p < \infty$.

Now consider the spaces $U_{\varepsilon}^{1,2} = W^{1,2}(\Omega) \oplus W^{1,2}(R_{\varepsilon})$ with the norm

$$||u_{\varepsilon}||_{U_{\varepsilon}^{1,2}}^{2} = ||u_{\varepsilon}||_{W^{1,2}(\Omega)}^{2} + \frac{1}{\varepsilon^{N-1}} ||u_{\varepsilon}||_{W^{1,2}(R_{\varepsilon})}^{2}$$
(2.8)

and $U_0^{1,2} = W^{1,2}(\Omega) \oplus W^{1,2}(0,1)$ with the norm

$$\|(w,v)\|_{U_0^{1,2}}^2 = \|w\|_{W^{1,2}(\Omega)}^2 + \int_0^1 g(|v_x|^2 + |v|^2).$$

Observe that the spaces $U_{\varepsilon}^{1,2}$ do not coincide algebraically with the spaces $W^{1,2}(\Omega_{\varepsilon})$ since we are allowing the functions of $U_{\varepsilon}^{1,2}$ to be discontinuous at $\partial\Omega \cap \partial R_{\varepsilon}$.

We also prove that

Theorem 2.7. The family of attractors $\{A_{\varepsilon}\}_{{\varepsilon}\in[0,1]}$ is E_{ε} -upper semicontinuous at ${\varepsilon}=0$ in $U_{\varepsilon}^{1,2}$.

Moreover, if every equilibria of the limit problem is hyperbolic, then the family of attractors is also E_{ε} - lower semicontinuous at $\varepsilon = 0$ in $U_{\varepsilon}^{1,2}$.

3. Convergence of resolvent operators

In this section we analyze the convergence of the resolvent operators associated to the elliptic operators A_{ε} defined in Section 2, that is, we study the convergence of $(A_{\varepsilon} + \lambda)^{-1} \to (A_0 + \lambda)^{-1}$ as $\varepsilon \to 0$ with λ in some region of the complex plane.

The convergence of resolvent operators is used, in Section 4, to analyze the convergence properties of the linear semigroups $e^{-A_{\varepsilon}t} \to e^{-A_0t}$ as $\varepsilon \to 0$, with the aid of the expression

$$e^{-A_{\varepsilon}t} = \frac{1}{2\pi i} \int_{\Gamma} e^{\lambda t} (A_{\varepsilon} + \lambda)^{-1} d\lambda, \quad t > 0$$

where Γ is an appropriate unbounded curve in the complex plane.

Moreover, since we need to analyze also the convergence properties of the linear semigroups associated to linearized equations around equilibria, that is $e^{-(A_{\varepsilon}-f'(u_{\varepsilon}^*))t}$ to $e^{-(A_0-f'(u_0^*))t}$ as ε

tends to 0, where u_{ε}^* and u_0^* are equilibria for (2.6), and (2.7), respectively, we will also need to study the convergence properties of the resolvent operators $(A_{\varepsilon}+V_{\varepsilon}+\lambda)^{-1} \to (A_0+V_0+\lambda)^{-1}$ as $\varepsilon \to 0$ for the potentials $V_{\varepsilon}(x) = -f'(u_{\varepsilon}^*(x))$ and $V_0(x) = -f'(u_0^*(x))$. To show this convergence we will need to obtain some rates of convergence of the equilibria u_{ε}^* to u_0^* .

We have divided the section in several subsections. In Subsection 3.1, we analyze the convergence of the resolvent operators for a fixed potential and in Subsection 3.2 we analyze the case of a potential which depends on the parameter ε . In Subsection 3.3 we obtain some rates of convergence of the equilibria and use this rates to obtain the convergence of the resolvent operators of the linearized operators around the equilibria.

3.1. Rate of convergence of resolvent operators: The case of a fixed potential. Consider a complex potential $V_0 = (V_{\Omega}, V_{R_0}) \in U_0^{\infty}$. Often, we write V_0 for $E_{\varepsilon}V_0 \in L^{\infty}(\Omega_{\varepsilon})$. Consider also the operator in $\mathcal{L}(L^p(\Omega_{\varepsilon}))$ and in $\mathcal{L}(U_0^p)$ which is the multiplication by the potential V_0 . We denote this operator again by V_0 , that is, $V_0(u_{\varepsilon}) \equiv (E_{\varepsilon}V_0)u_{\varepsilon} \equiv V_0u_{\varepsilon}$ and $V_0(w,v) = (V_{\Omega}w,V_{R_0}v)$.

Let us assume that $\operatorname{Re} \sigma(A_0 + V_0) > \delta > 0$. It follows from of [3, Proposition 3.13, Corollary 3.14] that, for all suitably small ε , $\operatorname{Re} \sigma(A_\varepsilon + V_0) \geqslant \delta > 0$.

The operator $A_{\varepsilon} + V_0$ is sectorial and the following estimate holds

$$\|(\lambda + A_{\varepsilon} + V_0)^{-1}\|_{\mathcal{L}(L^p(\Omega_{\varepsilon}))} \leqslant \frac{C}{|\lambda| + 1}, \quad \text{for} \quad \lambda \in \Sigma_{\theta},$$
 (3.1)

where $\Sigma_{\theta} = \{\lambda \in \mathbb{C} : |\arg(\lambda)| \leq \pi - \theta\}$, $0 < \theta < \frac{\pi}{2}$ and C is a constant that does not depend on ε , although it depends on p and blows up as $p \to \infty$. This estimate follows from the fact that the localization of the numerical range in the complex plane can be done independently of ε , see [13].

We know that, for any $0 < \varepsilon \le 1$, the operator $A_{\varepsilon} + V_0$ is a sectorial operator in U_{ε}^p and the following result holds

Lemma 3.1. For any bounded linear operator $J: L^p(\Omega_{\varepsilon}) \to L^p(\Omega_{\varepsilon})$ we have

$$||J||_{\mathcal{L}(U_{\varepsilon}^{p})} \leqslant ||J||_{\mathcal{L}(L^{p}(\Omega_{\varepsilon}), U_{\varepsilon}^{p})} \leqslant \varepsilon^{\frac{-N+1}{p}} ||J||_{\mathcal{L}(L^{p}(\Omega_{\varepsilon}))}$$
(3.2)

Proof: The proof of this result follows immediately from the norm estimate

$$\|\cdot\|_{U_{\varepsilon}^{p}} \leqslant \varepsilon^{\frac{-N+1}{p}} \|\cdot\|_{L^{p}(\Omega_{\varepsilon})}. \tag{3.3}$$

which follows directly from the definition of the norm in U_{ε}^{p} .

In particular, from Lemma 3.1 and from estimate (3.1), we have that for all $\lambda \in \Sigma_{\theta}$

$$\|(\lambda + A_{\varepsilon} + V_0)^{-1}\|_{\mathcal{L}(U_{\varepsilon}^p)} \leqslant \|(\lambda + A_{\varepsilon} + V_0)^{-1}\|_{\mathcal{L}(L^p(\Omega_{\varepsilon}), U_{\varepsilon}^p)} \leqslant C \frac{\varepsilon^{\frac{-N+1}{p}}}{|\lambda| + 1}, \quad \text{for} \quad \lambda \in \Sigma_{\theta}. \quad (3.4)$$

As for the limit problem, from [4], we have the following result.

Proposition 3.2. The operator $A_0 + V_0$ defined by (2.2) has the following properties

- i) $D(A_0 + V_0)$ is dense in U_0^p ,
- ii) $A_0 + V_0$ is a closed operator,
- iii) $A_0 + V_0$ has compact resolvent and

iv) $A_0 + V_0 : D(A_0 + V_0) \subset U_0^p \to U_0^p$ is such that, $\rho(A_0 + V_0) \supset \Sigma_\theta$ where $\Sigma_\theta = \{\lambda \in \mathbb{C} : |\operatorname{arg}(\lambda)| \leq \pi - \theta\}, \ 0 < \theta < \frac{\pi}{2}, \ and \ for \ p \geqslant q > \frac{N}{2},$

$$\|(\lambda + A_0 + V_0)^{-1}\|_{\mathcal{L}(U_0^q, U_0^p)} \le \frac{C}{|\lambda|^{\alpha} + 1},$$
 (3.5)

$$\|(\lambda + A_0 + V_0)^{-1}\|_{\mathcal{L}(U_0^{\infty})} \le \frac{C}{|\lambda| + 1},$$
 (3.6)

$$\|(\lambda + A_0 + V_0)^{-1}\|_{\mathcal{L}(U_0^{\infty}, U_0^p)} \le \frac{C}{|\lambda| + 1},$$
 (3.7)

for each $0 < \alpha < 1 - \frac{N}{2q} - \frac{1}{2}(\frac{1}{q} - \frac{1}{p}) < 1$ and $\lambda \in \Sigma_{\theta}$.

v) If B_0 is the realization of A_0 in $C(\overline{\Omega}) \oplus L^p(0,1)$ we have that B_0 is a sectorial operator in $C(\overline{\Omega}) \oplus L^p_g(0,1)$ with compact resolvent. Therefore $-B_0$ generates an analytic semigroup e^{-B_0t} in $C(\overline{\Omega}) \oplus L^p_g(0,1)$.

The following result is crucial to the remaining results in this section and to the whole program of the paper.

Proposition 3.3. If p > N and $2 \leq q < \infty$, there is a constant C, independent of ε , such that

$$||A_{\varepsilon}^{-1} f_{\varepsilon} - E_{\varepsilon} A_0^{-1} M_{\varepsilon} f_{\varepsilon}||_{H^1(\Omega) \oplus H^1(R_{\varepsilon})} \leqslant C \varepsilon^{N/2} ||f_{\varepsilon}||_{U_{\varepsilon}^p}, \tag{3.8}$$

$$||A_{\varepsilon}^{-1} f_{\varepsilon} - E_{\varepsilon} A_0^{-1} M_{\varepsilon} f_{\varepsilon}||_{L^{q}(\Omega_{\varepsilon})} \leq C \varepsilon^{N/q} ||f_{\varepsilon}||_{U_{\varepsilon}^{p}}, \tag{3.9}$$

and

$$||A_{\varepsilon}^{-1} f_{\varepsilon} - E_{\varepsilon} A_{0}^{-1} M_{\varepsilon} f_{\varepsilon}||_{U_{\varepsilon}^{q}} \leqslant C \varepsilon^{1/q} ||f_{\varepsilon}||_{U_{\varepsilon}^{p}}, \tag{3.10}$$

for all $f_{\varepsilon} \in U_{\varepsilon}^p$.

Proof: The inequality (3.8) was proved in Proposition A. 8 in [3]. This estimate is the key estimate for [3] and also for the complete analysis we are performing in the dumbbell domains.

Observe that in particular, from (3.8), we obtain that

$$||A_{\varepsilon}^{-1} f_{\varepsilon} - E_{\varepsilon} A_{0}^{-1} M_{\varepsilon} f_{\varepsilon}||_{L^{2}(\Omega_{\varepsilon})} \leqslant C \varepsilon^{N/2} ||f_{\varepsilon}||_{U_{\varepsilon}^{p}}.$$
(3.11)

From [3, Lemma A.11], for p > N/2 we have

$$||A_{\varepsilon}^{-1} f_{\varepsilon}||_{L^{\infty}(\Omega_{\varepsilon})} \leqslant C ||f_{\varepsilon}||_{U_{\varepsilon}^{p}}. \tag{3.12}$$

Also we know that if p > N/2, $||A_0^{-1}M_{\varepsilon}f_{\varepsilon}||_{L^{\infty}(\Omega) \oplus L^{\infty}(0,1)} \leqslant C||M_{\varepsilon}f_{\varepsilon}||_{L^{p}(\Omega) \oplus L^{p}(0,1)}$ then

$$||E_{\varepsilon}A_0^{-1}M_{\varepsilon}f_{\varepsilon}||_{L^{\infty}(\Omega_{\varepsilon})} \leqslant C||f_{\varepsilon}||_{U_{\varepsilon}^{p}}.$$
(3.13)

which implies that

$$||A_{\varepsilon}^{-1} f_{\varepsilon} - E_{\varepsilon} A_0^{-1} M_{\varepsilon} f_{\varepsilon}||_{L^{\infty}(\Omega_{\varepsilon})} \leqslant C ||f_{\varepsilon}||_{U_{\varepsilon}^{p}}.$$
(3.14)

For $q \ge 2$, (3.9) follows from (3.11) and (3.14) and interpolation. The estimate (3.10) follows from (3.9) and (3.3).

To obtain the resolvent convergence of $A_{\varepsilon} + V_0$ we strongly use the previous result and the following uniform (with respect to ε) estimate.

Lemma 3.4. If V_0 is such that $(A_0 + V_0)$ is invertible, for $p > \frac{N}{2}$, we have

$$||E_{\varepsilon}(A_0 + V_0)^{-1} M_{\varepsilon}||_{\mathcal{L}(U_{\varepsilon}^p)} \leqslant C \tag{3.15}$$

and, for each $p > \frac{N}{2}$, there is a constant C, independent of ε , such that

$$||E_{\varepsilon}(A_0 + V_0)^{-1} M_{\varepsilon}||_{\mathcal{L}(L^p(\Omega_{\varepsilon}))} \leqslant C.$$
(3.16)

Proof: Statement (3.15) follows from $||E_{\varepsilon}||_{\mathcal{L}(U_0^p,U_{\varepsilon}^p)} = ||M_{\varepsilon}||_{\mathcal{L}(U_{\varepsilon}^p,U_0^p)} = 1$ (see [3]) and from Proposition 3.2.

For (3.16) we proceed as follows. Let $f_{\varepsilon} \in L^p(\Omega_{\varepsilon})$ and $u_{\varepsilon} = (w_{\varepsilon}, v_{\varepsilon}) = (A_0 + V_0)^{-1} M_{\varepsilon} f_{\varepsilon}$, then

$$\begin{cases}
-\Delta w_{\varepsilon} + w_{\varepsilon} + V_{\Omega}(x)w_{\varepsilon} = f_{\varepsilon}, & \Omega, \\
\frac{\partial w_{\varepsilon}}{\partial n} = 0, & \partial\Omega \\
-\frac{1}{g} (g(v_{\varepsilon})_{s})_{s} + v_{\varepsilon} + V_{R_{0}}(s)v_{\varepsilon} = M_{\varepsilon}f_{\varepsilon}, & (0, 1) \\
v_{\varepsilon}(0) = w_{\varepsilon}(P_{0}), & v_{\varepsilon}(1) = w_{\varepsilon}(P_{1}).
\end{cases}$$

Since $p > \frac{N}{2}$, we have that

$$||w_{\varepsilon}||_{L^{p}(\Omega)} \leqslant C ||f_{\varepsilon}||_{L^{p}(\Omega)} \quad \text{and} \quad ||w_{\varepsilon}||_{C(\bar{\Omega})} \leqslant C ||f_{\varepsilon}||_{L^{p}(\Omega)}.$$

In particular $|w_{\varepsilon}(P_0)| + |w_{\varepsilon}(P_1)| \leq C ||f_{\varepsilon}||_{L^p(\Omega)}$. Also

$$||v_{\varepsilon}||_{L^{p}(0,1)} \leq |w_{\varepsilon}(P_{0})| + |w_{\varepsilon}(P_{1})| + ||M_{\varepsilon}f_{\varepsilon}||_{L^{p}(0,1)}$$

and

$$||E_{\varepsilon}v_{\varepsilon}||_{L^{p}(R_{\varepsilon})} = \varepsilon^{\frac{N-1}{p}} ||v_{\varepsilon}||_{L^{p}(0,1)} \leqslant \varepsilon^{\frac{N-1}{p}} (|w_{\varepsilon}(P_{0})| + |w_{\varepsilon}(P_{1})|) + \varepsilon^{\frac{N-1}{p}} ||M_{\varepsilon}f_{\varepsilon}||_{L^{p}(0,1)}$$
$$\leqslant |w_{\varepsilon}(P_{0})| + |w_{\varepsilon}(P_{1})| + ||f_{\varepsilon}||_{L^{p}(R_{\varepsilon})}$$
$$\leqslant C ||f_{\varepsilon}||_{L^{p}(\Omega_{\varepsilon})}.$$

where we have used that $||M_{\varepsilon}f_{\varepsilon}||_{L^{p}(0,1)} \leqslant \varepsilon^{-\frac{N-1}{p}} ||f_{\varepsilon}||_{L^{p}(R_{\varepsilon})}$. The proof is now complete.

The next two lemmas are resolvent identities which allow us (together with the previous lemma) to transfer information from the resolvent convergence of A_{ε} to the resolvent convergence of $A_{\varepsilon} + V_0$.

Lemma 3.5. If $(A_0 + V_0)$ and $(A_{\varepsilon} + V_0)$ are both invertible the following identity holds

$$(A_{\varepsilon} + V_0)^{-1} - E_{\varepsilon} (A_0 + V_0)^{-1} M_{\varepsilon}$$

$$= [I - (A_{\varepsilon} + V_0)^{-1} V_0] (A_{\varepsilon}^{-1} - E_{\varepsilon} A_0^{-1} M_{\varepsilon}) [I - E_{\varepsilon} V_0 (A_0 + V_0)^{-1} M_{\varepsilon}].$$
(3.17)

Proof: Since $(I - (A_{\varepsilon} + V_0)^{-1}V_0)(I + A_{\varepsilon}^{-1}V_0) = I$, the identity (3.17) is equivalent to

$$(A_{\varepsilon}^{-1} - E_{\varepsilon} A_0^{-1} M_{\varepsilon}) (I - E_{\varepsilon} V_0 (A_0 + V_0)^{-1} M_{\varepsilon}) = = (I + A_{\varepsilon}^{-1} V_0) ((A_{\varepsilon} + V_0)^{-1} - E_{\varepsilon} (A_0 + V_0)^{-1} M_{\varepsilon}).$$
(3.18)

Using that $V_0(A_0 + V_0)^{-1} = I - A_0(A_0 + V_0)^{-1}$ and expanding the left hand side of (3.18) we have

$$(A_{\varepsilon}^{-1} - E_{\varepsilon} A_0^{-1} M_{\varepsilon}) (I - E_{\varepsilon} V_0 (A_0 + V_0)^{-1} M_{\varepsilon}) = A_{\varepsilon}^{-1} - A_{\varepsilon}^{-1} E_{\varepsilon} V_0 (A_0 + V_0)^{-1} M_{\varepsilon}$$
$$- E_{\varepsilon} A_0^{-1} M_{\varepsilon} + E_{\varepsilon} A_0^{-1} (I - A_0 (A_0 + V_0)^{-1}) M_{\varepsilon}$$
$$= A_{\varepsilon}^{-1} - A_{\varepsilon}^{-1} E_{\varepsilon} V_0 (A_0 + V_0)^{-1} M_{\varepsilon} - E_{\varepsilon} (A_0 + V_0)^{-1} M_{\varepsilon}$$

On the other hand, using that $A_{\varepsilon}^{-1} = (I + A_{\varepsilon}^{-1}V_0)(A_{\varepsilon} + V_0)^{-1}$ and expanding the right hand side of (3.18), we have

$$(I + A_{\varepsilon}^{-1}V_0)((A_{\varepsilon} + V_0)^{-1} - E_{\varepsilon}(A_0 + V_0)^{-1}M_{\varepsilon})$$

= $A_{\varepsilon}^{-1} - E_{\varepsilon}(A_0 + V_0)^{-1}M_{\varepsilon} - A_{\varepsilon}^{-1}E_{\varepsilon}V_0(A_0 + V_0)^{-1}M_{\varepsilon}.$

which proves (3.18).

In a very similar way we also have,

Lemma 3.6. If $(A_0 + V_0)$ and $(A_{\varepsilon} + V_0)$ are both invertible, the following identity holds

$$(A_{\varepsilon} + V_0)^{-1} - E_{\varepsilon}(A_0 + V_0)^{-1} M_{\varepsilon} = [I - E_{\varepsilon}(A_0 + V_0)^{-1} V_0 M_{\varepsilon}] (A_{\varepsilon}^{-1} - E_{\varepsilon} A_0^{-1} M_{\varepsilon}) \quad [I - V_0 (A_{\varepsilon} + V_0)^{-1}].$$
(3.19)

Proof: The proof is similar to the one provided for the previous lemma.

We are now ready to prove the main results of this section

Proposition 3.7. If p,q > N, $(A_0 + V_0) : D(A_0) \subset U_0^p \to U_0^p$ has bounded inverse and $f_{\varepsilon} \in U_{\varepsilon}^p$, then

$$\|(A_{\varepsilon} + V_0)^{-1} f_{\varepsilon} - E_{\varepsilon} (A_0 + V_0)^{-1} M_{\varepsilon} f_{\varepsilon} \|_{L^q(\Omega_{\varepsilon})} \leqslant C \varepsilon^{N/q} \|f_{\varepsilon}\|_{U_{\varepsilon}^p}, \tag{3.20}$$

where C depends on $\|(A_0+V_0)^{-1}\|_{\mathcal{L}(U_0^p,U_0^p)}$ and on $\|V_0\|_{L^{\infty}}$, but not on ε or f_{ε} .

Proof: Let us start pointing out that if $(A_0 + V_0)$ is invertible, from [3] we also have that $(A_{\varepsilon} + V_0)$ is invertible for all suitably small ε . Hence (3.20) makes sense.

Adding and subtracting the appropriate term in (3.17) we have:

$$(A_{\varepsilon} + V_{0})^{-1} - E_{\varepsilon}(A_{0} + V_{0})^{-1} M_{\varepsilon}$$

$$= (-(A_{\varepsilon} + V_{0})^{-1} + E_{\varepsilon}(A_{0} + V_{0})^{-1} M_{\varepsilon}) V_{0} (A_{\varepsilon}^{-1} - E_{\varepsilon} A_{0}^{-1} M_{\varepsilon}) (I - V_{0} E_{\varepsilon} (A_{0} + V_{0})^{-1} M_{\varepsilon})$$

$$+ (I - E_{\varepsilon} (A_{0} + V_{0})^{-1} M_{\varepsilon} V_{0}) (A_{\varepsilon}^{-1} - E_{\varepsilon} A_{0}^{-1} M_{\varepsilon}) (I - V_{0} E_{\varepsilon} (A_{0} + V_{0})^{-1} M_{\varepsilon}).$$

Let us first estimate

$$\Theta_{\varepsilon} = ((A_{\varepsilon} + V_0)^{-1} - E_{\varepsilon}(A_0 + V_0)^{-1}M_{\varepsilon})V_0(A_{\varepsilon}^{-1} - E_{\varepsilon}A_0^{-1}M_{\varepsilon})(I - V_0E_{\varepsilon}(A_0 + V_0)^{-1}M_{\varepsilon})$$

Note that, from inequality (3.10) and (3.9) we have that

$$\|A_{\varepsilon}^{-1} - E_{\varepsilon}A_0^{-1}M_{\varepsilon}\|_{\mathcal{L}(U_{\varepsilon}^p,U_{\varepsilon}^p)} \leqslant C\varepsilon^{1/p} \text{ and } \|A_{\varepsilon}^{-1} - E_{\varepsilon}A_0^{-1}M_{\varepsilon}\|_{\mathcal{L}(U_{\varepsilon}^p,L^q(\Omega_{\varepsilon}))} \leqslant C\varepsilon^{N/q}.$$

Since

$$||V_0||_{\mathcal{L}(L^q(\Omega_{\varepsilon}))} \leqslant C||V_0||_{L^{\infty}(\Omega_{\varepsilon})} \text{ and } ||V_0||_{\mathcal{L}(U^p_{\varepsilon})} \leqslant C||V_0||_{L^{\infty}(\Omega_{\varepsilon})},$$

it follows from (3.15) that

$$\|\Theta_{\varepsilon}\|_{\mathcal{L}(U_{\varepsilon}^{p},L^{q}(\Omega_{\varepsilon}))} \leqslant C\varepsilon^{1/p}\|(A_{\varepsilon}+V_{0})^{-1}-E_{\varepsilon}(A_{0}+V_{0})^{-1}M_{\varepsilon}\|_{\mathcal{L}(U_{\varepsilon}^{p},L^{q}(\Omega_{\varepsilon}))}.$$

where $C = C(\|V_0\|_{L^{\infty}(\Omega_{\varepsilon})})$ is independent of ε . Choosing ε_0 such that $C\varepsilon^{1/p} \leqslant \frac{1}{2}$, for all $\varepsilon \in [0, \varepsilon_0]$, we have that

$$\|(A_{\varepsilon}+V_0)^{-1}-E_{\varepsilon}(A_0+V_0)^{-1}M_{\varepsilon}\|_{\mathcal{L}(U_{\varepsilon}^p,L^q(\Omega_{\varepsilon}))} \leq 2\|(I-E_{\varepsilon}(A_0+V_0)^{-1}M_{\varepsilon}V_0)(A_{\varepsilon}^{-1}-E_{\varepsilon}A_0^{-1}M_{\varepsilon})(I-V_0E_{\varepsilon}(A_0+V_0)^{-1}M_{\varepsilon})\|_{\mathcal{L}(U_{\varepsilon}^p,L^q(\Omega_{\varepsilon}))}.$$

Now, from (3.15) and (3.16) there is a constant C, independent of ε , such that

$$||(I - V_0 E_{\varepsilon} (A_0 + V_0)^{-1} M_{\varepsilon})||_{\mathcal{L}(U_{\varepsilon}^p)} \leqslant 1 + C ||V_0||_{L^{\infty}(\Omega_{\varepsilon})},$$

$$||(I - E_{\varepsilon} (A_0 + V_0)^{-1} M_{\varepsilon} V_0)||_{\mathcal{L}(L^q(\Omega_{\varepsilon}))} \leqslant 1 + C ||V_0||_{L^{\infty}(\Omega_{\varepsilon})},$$

Therefore, using (3.9),

$$\|(A_{\varepsilon}+V_0)^{-1}-E_{\varepsilon}(A_0+V_0)^{-1}M_{\varepsilon}\|_{\mathcal{L}(U_{\varepsilon}^p,L^q(\Omega_{\varepsilon}))} \leqslant C \varepsilon^{N/q},$$

where the constant C does depends on $||V_0||_{L^{\infty}(\Omega_{\varepsilon})}$. This shows the proposition.

3.2. Rate of convergence of resolvent operators: The case of a varying potential. We are going to study now the convergence properties of resolvent operators of the form $(A_{\varepsilon} + W_{\varepsilon})^{-1}$ to $(A_0 + W_0)^{-1}$, where W_{ε} converges to W_0 in a sense to be specified. We need to perform this study since we want to compare the resolvent operators of the linearizations around equilibria. Hence, we will have a family of equilibria u_{ε}^* which will converge to an equilibria of the limiting problem u_0^* and we will need to consider the operators $A_{\varepsilon} - f'(u_{\varepsilon}^*)$ and $A_0 - f'(u_0^*)$ and analyze the convergence properties of their resolvent.

Having this in mind, let us consider the following setting for the potentials,

(H) $V_{\varepsilon} \in L^{\infty}(\Omega_{\varepsilon})$, $V_0 = (V_{\Omega}, V_{R_0}) \in U_0^{\infty}$ be two potentials which satisfy that $|V_{\varepsilon}|, |V_0| \leq a$ for some a > 0 and such that for $N < q < \infty$ we have

$$\varepsilon^{\frac{-N+1}{q}} \| V_{\varepsilon} - E_{\varepsilon} V_0 \|_{L^q(\Omega_{\varepsilon})} \to 0, \text{ as } \varepsilon \to 0$$
 (3.21)

Denote by $W_{\varepsilon} = V_{\varepsilon} + a$, $W_0 = V_0 + a = (V_{\Omega} + a, V_{R_0} + a)$ so that W_{ε} and W_0 are positive and they also satisfy an estimate like (3.21) substituting V_{ε} and V_0 by W_{ε} and W_0 respectively.

As we did in Subsection 3.1, let us identify the potentials W_{ε} , W_0 with their corresponding multiplication operators.

With this notation and writing $\Lambda_{\varepsilon} = A_{\varepsilon} + W_{\varepsilon}$, we have that the operator Λ_{ε} is sectorial and the following estimate holds

$$\|(\lambda + \Lambda_{\varepsilon})^{-1}\|_{\mathcal{L}(L^{p}(\Omega_{\varepsilon}))} \leqslant \frac{C}{|\lambda| + 1}, \quad \text{for} \quad \lambda \in \Sigma_{\theta},$$
 (3.22)

where $\Sigma_{\theta} = \{\lambda \in \mathbb{C} : |\arg(\lambda)| \leq \pi - \theta\}$, $0 < \theta < \frac{\pi}{2}$ and C is a constant that does not depend on ε (that follows form the fact that the localization of the numerical range in the complex plane can be done independently of ε), however it depends on p and blows up as $p \to \infty$, see [13].

We know that, for any $0 < \varepsilon \leq 1$, the operator Λ_{ε} is a sectorial operator in U_{ε}^{p} and the following result holds

Lemma 3.8. For all $\lambda \in \Sigma_{\theta}$ we have that

$$\|(\lambda + \Lambda_{\varepsilon})^{-1}\|_{\mathcal{L}(U_{\varepsilon}^{p})} \leqslant \|(\lambda + \Lambda_{\varepsilon})^{-1}\|_{\mathcal{L}(L^{p}(\Omega_{\varepsilon}), U_{\varepsilon}^{p})} \leqslant C \frac{\varepsilon^{\frac{-N+1}{p}}}{|\lambda| + 1}.$$
 (3.23)

Proof: It follows immediately from (3.22) and from Lemma 3.1.

The following result follows easily from the properties of resolvent operators. It is crucial to obtain convergence properties for resolvent operators from the convergence properties of $\Lambda_{\varepsilon}^{-1}$ to Λ_{0}^{-1} .

Lemma 3.9. As an immediate consequence of (3.5), (3.6) and (3.7), there is a constant C such that, for all $\lambda \in \Sigma_{\theta}$, $p \geqslant q > \frac{N}{2}$ and $0 < \alpha < 1 - \frac{N}{2q} - \frac{1}{2}(\frac{1}{q} - \frac{1}{p}) < 1$

$$||E_{\varepsilon}(\lambda + \Lambda_0)^{-1} M_{\varepsilon}||_{\mathcal{L}(U_{\varepsilon}^q, U_{\varepsilon}^p)} \leqslant \frac{C}{|\lambda|^{\alpha} + 1}, \tag{3.24}$$

$$||E_{\varepsilon}(\lambda + \Lambda_0)^{-1}M_{\varepsilon}||_{\mathcal{L}(C(\bar{\Omega}_{\varepsilon}), L^{\infty}(\Omega_{\varepsilon}))} \leq \frac{C}{|\lambda| + 1},$$
 (3.25)

and

$$||E_{\varepsilon}\lambda(\lambda+\Lambda_0)^{-1}M_{\varepsilon}||_{\mathcal{L}(C(\bar{\Omega}_{\varepsilon}),U_{\varepsilon}^p)} \leqslant C$$
(3.26)

where C is a constant that does not depend in ε .

We have now the following key result, which is analogous to Proposition 3.3 and Proposition 3.7

Proposition 3.10. For p, q > N and $f_{\varepsilon} \in U_{\varepsilon}^{p}$ we have

$$\|\Lambda_{\varepsilon}^{-1} f_{\varepsilon} - E_{\varepsilon} \Lambda_{0}^{-1} M_{\varepsilon} f_{\varepsilon}\|_{L^{q}(\Omega_{\varepsilon})} \leqslant C(\varepsilon^{\frac{N}{q}} + \|W_{\varepsilon} - E_{\varepsilon} W_{0} M_{\varepsilon}\|_{L^{q}(\Omega_{\varepsilon})}) \|f_{\varepsilon}\|_{U_{\varepsilon}^{p}}.$$

$$(3.27)$$

with C independent of ε and f_{ε} .

Proof: Let $f_{\varepsilon} \in U_{\varepsilon}^{p}$ and let $u_{\varepsilon} = \Lambda_{\varepsilon}^{-1} f_{\varepsilon} = (A_{\varepsilon} + W_{\varepsilon})^{-1} f_{\varepsilon}$. Consider the auxiliary function, $\tilde{u}_{\varepsilon} = (A_{\varepsilon} + E_{\varepsilon} W_{0})^{-1} f_{\varepsilon}$, i.e.,

$$\begin{cases}
-\Delta u_{\varepsilon} + u_{\varepsilon} + W_{\varepsilon}u_{\varepsilon} &= f_{\varepsilon}, \quad \Omega_{\varepsilon}, \\
\frac{\partial u_{\varepsilon}}{\partial n} &= 0, \quad \partial \Omega_{\varepsilon}.
\end{cases}$$
(3.28)

$$\begin{cases}
-\Delta \tilde{u}_{\varepsilon} + \tilde{u}_{\varepsilon} + W_0 \tilde{u}_{\varepsilon} &= f_{\varepsilon}, \quad \Omega_{\varepsilon}, \\
\frac{\partial \tilde{u}_{\varepsilon}}{\partial n} &= 0, \quad \partial \Omega_{\varepsilon}.
\end{cases}$$
(3.29)

From comparison results, it is easy to see that $|\tilde{u}_{\varepsilon}| \leq \bar{u}_{\varepsilon}$ where

$$\begin{cases} -\Delta \bar{u}_{\varepsilon} + \bar{u}_{\varepsilon} &= |f_{\varepsilon}|, \quad \Omega_{\varepsilon}, \\ \frac{\partial u_{\varepsilon}}{\partial n} &= 0, \quad \partial \Omega_{\varepsilon}. \end{cases}$$

Applying Lemma A.11 of [3], we have that

$$\|\bar{u}_{\varepsilon}\|_{L^{\infty}(\Omega_{\varepsilon})} \leqslant C \|f_{\varepsilon}\|_{U_{\varepsilon}^{p}} \quad \text{for } p > N/2$$
 (3.30)

which implies

$$\|\tilde{u}_{\varepsilon}\|_{L^{\infty}(\Omega_{\varepsilon})} \leqslant C \|f_{\varepsilon}\|_{U_{\varepsilon}^{p}}.$$

Next, observe that

$$u_{\varepsilon} = (A_{\varepsilon} + E_{\varepsilon}W_0)^{-1} f_{\varepsilon} + (A_{\varepsilon} + E_{\varepsilon}W_0)^{-1} (E_{\varepsilon}W_0 - W_{\varepsilon}) u_{\varepsilon}$$
$$u_0 = (A_0 + W_0)^{-1} M_{\varepsilon} f_{\varepsilon}.$$

Hence,

$$\|u_{\varepsilon} - E_{\varepsilon}u_{0}\|_{L^{q}(\Omega_{\varepsilon})}$$

$$\leq \|(A_{\varepsilon} + E_{\varepsilon}W_{0})^{-1} - E_{\varepsilon}(A_{0} + W_{0})^{-1}M_{\varepsilon}f_{\varepsilon}\|_{L^{q}(\Omega_{\varepsilon})} + \|(A_{\varepsilon} + E_{\varepsilon}W_{0})^{-1}(W_{\varepsilon} - E_{\varepsilon}W_{0})u_{\varepsilon}\|_{L^{q}(\Omega_{\varepsilon})}$$

$$\leq C\varepsilon^{\frac{N}{q}}\|f_{\varepsilon}\|_{U_{\varepsilon}^{p}} + \bar{C}\|(A_{\varepsilon} + E_{\varepsilon}W_{0})^{-1}\|_{\mathcal{L}(L^{q}(\Omega_{\varepsilon}))}\|W_{\varepsilon} - E_{\varepsilon}W_{0}\|_{L^{q}(\Omega_{\varepsilon})}\|u_{\varepsilon}\|_{L^{\infty}(\Omega_{\varepsilon})}$$

$$\leq \tilde{C}(\varepsilon^{\frac{N}{q}} + \|W_{\varepsilon} - E_{\varepsilon}W_{0}\|_{L^{q}(\Omega_{\varepsilon})})\|f_{\varepsilon}\|_{U^{p}}.$$

where we have used (3.20) and the fact that that there is a constant C, independent of ε and of $q \in [1, \infty]$, such that $\|(A_{\varepsilon} + W_0)^{-1}\|_{\mathcal{L}(L^q(\Omega_{\varepsilon}))} \leq C$. This shows the lemma.

As an immediate corollary, we have

Corollary 3.11. For p, q > N we have

$$\varepsilon^{-\frac{N-1}{q}} \|\Lambda_{\varepsilon}^{-1} - E_{\varepsilon} \Lambda_0^{-1} M_{\varepsilon}\|_{\mathcal{L}(U_{\varepsilon}^p, L^q(\Omega_{\varepsilon}))} \to 0 \text{ as } \varepsilon \to 0.$$
 (3.31)

Proof: We just need to apply the previous proposition and hypothesis (H).

Now consider a compact subset K of the complex plane which is contained in the resolvent set of the operator Λ_0 . Let c(K) be a positive constant such that

$$\sup_{\lambda \in K} \|(\lambda + \Lambda_0)^{-1}\|_{\mathcal{L}(U_0^p, U_0^p)} \leqslant c(K).$$

Also, let $\Sigma_{\theta} := \{z \in C : |\arg(z)| \leq \pi - \theta\}$, for $0 < \theta < \pi/2$.

Proposition 3.12. For p,q > N, there exists a constant $C = C(K,\theta)$, a number $\varepsilon_0 > 0$ and a function $\eta(\varepsilon) \to 0$ as $\varepsilon \to 0$ such that for each $\lambda \in K \cup \Sigma_{\theta}$ and $0 < \varepsilon \leqslant \varepsilon_0$ we have

$$\varepsilon^{-\frac{N-1}{q}} \| (\lambda + \Lambda_{\varepsilon})^{-1} - E_{\varepsilon} (\lambda + \Lambda_{0})^{-1} M_{\varepsilon} \|_{\mathcal{L}(U_{\varepsilon}^{p}, L^{q}(\Omega_{\varepsilon}))} \leqslant C (1 + |\lambda|^{1-\alpha}) \eta(\varepsilon), \tag{3.32}$$

where $0 < \alpha < 1 - \frac{N}{2n} < 1$

Proof: Observe first that the spectrum of the operators Λ_{ε} and Λ_0 are subsets of $[1, +\infty)$.

Hence, if $\lambda \in \Sigma_{\theta}$ both $(\lambda + \Lambda_{\varepsilon})^{-1}$ and $(\lambda + \Lambda_{0})^{-1}$ make perfect sense for $0 < \varepsilon \leqslant \varepsilon_{0}$. Moreover, by the compact convergence of $A_{\varepsilon}^{-1} \to A_{0}^{-1}$, the convergence of $W_{\varepsilon} \to W_{0}$ and since $\|(\lambda + \Lambda_0)^{-1}\|_{\mathcal{L}(U_0^p, U_0^p)} = \|(\lambda + V_0 + A_0)^{-1}\|_{\mathcal{L}(U_0^p, U_0^p)} \leqslant c(K)$ for each $\lambda \in K$ which is a compact set in C, we have that $(\lambda + A_{\varepsilon} + V_{\varepsilon})$ and $(\lambda + A_{\varepsilon} + V_{0})$ are invertible for $0 < \varepsilon < \varepsilon_{0}$ and $\lambda \in \Lambda_0$ and $\|(\lambda + \Lambda_{\varepsilon})^{-1}\|_{\mathcal{L}(U_0^p, U_0^p)} \leq \tilde{c}(K)$, for some constant $\tilde{c}(K)$ and for all $\lambda \in K$. If this is not the case, then we could get a sequence of $\varepsilon_n \to 0$ and $\lambda_n \to \tilde{\lambda} \in K$ such that $\|(\lambda_n + \Lambda_{\varepsilon_n})^{-1}\|_{\mathcal{L}(U_0^p, U_0^p)} \to +\infty$. But this is in contradiction with the compact convergence of $(\lambda_n + \Lambda_{\varepsilon_n})^{-1}$ to $(\tilde{\lambda} + \Lambda_0)^{-1}$, see Lemma 4.7 of [3].

Hence, with this argument and with (3.22) and (3.24) we obtain

$$\|\lambda (\lambda + \Lambda_{\varepsilon})^{-1}\|_{\mathcal{L}(L^{q}(\Omega_{\varepsilon}))} \leqslant C, \quad \text{for} \quad \lambda \in K \cup \Sigma_{\theta},$$
 (3.33)

$$||E_{\varepsilon}\lambda(\lambda+\Lambda_0)^{-1}M_{\varepsilon}||_{\mathcal{L}(U_{\varepsilon}^p,U_{\varepsilon}^p)} \leqslant C(1+|\lambda|^{1-\alpha}), \quad \text{for} \quad \lambda \in K \cup \Sigma_{\theta}.$$
 (3.34)

with $0 < \alpha < 1 - \frac{N}{2p} < 1$. Applying Lemma 3.5 with Λ_0 in place of A_0 and λ in place of V_0 , we have

$$\begin{split} &\|(\lambda + \Lambda_{\varepsilon})^{-1} - E_{\varepsilon}(\lambda + \Lambda_{0})^{-1} M_{\varepsilon} \|_{\mathcal{L}(U_{\varepsilon}^{p}, L^{q}(\Omega_{\varepsilon}))} \leqslant \\ &\|I + \lambda(\lambda + \Lambda_{\varepsilon})^{-1} \|_{\mathcal{L}(L^{q}(\Omega_{\varepsilon}))} \|\Lambda_{\varepsilon}^{-1} - E_{\varepsilon} \Lambda_{0}^{-1} M_{\varepsilon} \|_{\mathcal{L}(U_{\varepsilon}^{p}, L^{q}(\Omega_{\varepsilon}))} \|[I - E_{\varepsilon} \lambda(\lambda + \Lambda_{0})^{-1} M_{\varepsilon}] \|_{\mathcal{L}(U_{\varepsilon}^{p})} \\ &\leqslant C(1 + |\lambda|^{1-\alpha}) \|\Lambda_{\varepsilon}^{-1} - E_{\varepsilon} \Lambda_{0}^{-1} M_{\varepsilon} \|_{\mathcal{L}(U_{\varepsilon}^{p}, L^{q}(\Omega_{\varepsilon}))} \leqslant C \varepsilon^{\frac{N-1}{q}} (1 + |\lambda|^{1-\alpha}) \eta(\varepsilon) \end{split}$$

where $\eta(\varepsilon) = \varepsilon^{-\frac{N-1}{q}} \|\Lambda_{\varepsilon}^{-1} - E_{\varepsilon}\Lambda_0^{-1}M_{\varepsilon}\|_{\mathcal{L}(U_{\varepsilon}^p, L^q(\Omega_{\varepsilon}))} \to 0 \text{ as } \varepsilon \to 0 \text{ by Corollary 3.11. This proves the proposition.}$

Remark 3.13. The results of Proposition 3.12 also hold for the operator A_{ε} instead of Λ_{ε} , that is with $W_{\varepsilon} = W_0 = 0$.

Corollary 3.14. In the conditions of Proposition 3.12, we have the following estimates,

$$\|(\lambda + \Lambda_{\varepsilon})^{-1} - E_{\varepsilon}(\lambda + \Lambda_{0})^{-1} M_{\varepsilon}\|_{\mathcal{L}(U_{\varepsilon}^{p}, U_{\varepsilon}^{q})} \leq C (1 + |\lambda|^{1-\alpha}) \eta(\varepsilon), \tag{3.35}$$

$$\|(\lambda + \Lambda_{\varepsilon})^{-1}\|_{\mathcal{L}(U_{\varepsilon}^{p}, U_{\varepsilon}^{q})} \leq C(1 + |\lambda|^{1-\alpha})$$
(3.36)

Proof: To prove (3.35) we just use that $\varepsilon^{-\frac{N-1}{q}} \| \cdot \|_{\mathcal{L}(U_{\varepsilon}^p, L^q(\Omega_{\varepsilon}))} \leq \| \cdot \|_{\mathcal{L}(U_{\varepsilon}^p, U_{\varepsilon}^q)}$ in (3.32). To prove (3.36) we just use (3.35) and (3.24), to obtain

$$\|(\lambda + \Lambda_{\varepsilon})^{-1}\|_{\mathcal{L}(U_{\varepsilon}^p, U_{\varepsilon}^q)} \leqslant C(1 + |\lambda|^{1-\alpha})\eta(\varepsilon) + \frac{C}{|\lambda|^{\alpha} + 1} \leqslant C(1 + |\lambda|^{1-\alpha}),$$

as we wanted to show.

These results play a fundamental role on the convergence of the linear semigroups for it will ensure the uniform convergence of the integrals defining them and will allow us to pass to the limit.

3.3. Rate of convergence of hyperbolic equilibria and of its linearizations. In this Subsection we will obtain rates of convergence of hyperbolic equilibria which, besides being interesting by themselves, they show that if we consider the potentials $V_{\varepsilon} = -f'(u_{\varepsilon}^*)$, $V_0 = -f'(u_0^*)$ then hypothesis (**H**) from Subsection 3.2 is satisfied, with $a = \sup\{|f'(s)| : s \in \mathbb{R}\}$. This will ensure that all the results from Subsection 3.2 apply for $\Lambda_{\varepsilon} = A_{\varepsilon} - f'(u_{\varepsilon}^*) + a$ and $\Lambda_0 = A_0 - f'(u_0^*) + a$.

Proposition 3.15. Let u_0^* be a hyperbolic equilibrium for (1.2) and (from the results in [3]) let u_{ε}^* be the sequence of hyperbolic equilibria for (1.1) satisfying that u_{ε}^* E-converges to u_0^* . Then, for q > N, we have

$$||u_{\varepsilon}^* - E_{\varepsilon} u_0^*||_{L^q(\Omega)} \leqslant C \varepsilon^{\frac{N}{q}}$$
(3.37)

and

$$\varepsilon^{-\frac{N-1}{q}} \| u_{\varepsilon}^* - E_{\varepsilon} u_0^* \|_{U_{\varepsilon}^p} \to 0, \quad as \ \varepsilon \to 0.$$
 (3.38)

Proof: Let $u_0^* = (w_0^*, v_0^*)$ be a hyperbolic equilibrium point for (1.2) and u_{ε}^* an equilibrium point for (1.1) with $\|u_{\varepsilon}^* - E_{\varepsilon}u_0^*\|_{U_{\varepsilon}^p} \xrightarrow{\varepsilon \to 0} 0$. For $V_0(x) = -f'(u_0^*(x))$, we write

$$u_{\varepsilon}^* = (A_{\varepsilon} + V_0)^{-1} (f(u_{\varepsilon}^*) + V_0 u_{\varepsilon}^*)$$
 and $u_0^* = (A_0 + V_0)^{-1} (f(u_0^*) + V_0 u_0^*)$.

Hence, taking norms in $L^q(\Omega)$, we get

$$||u_{\varepsilon}^{*} - E_{\varepsilon}u_{0}^{*}||_{L^{q}(\Omega)} = ||(A_{\varepsilon} + V_{0})^{-1}(f(u_{\varepsilon}^{*}) + V_{0}u_{\varepsilon}^{*}) - E_{\varepsilon}(A_{0} + V_{0})^{-1}(f(u_{0}^{*}) + V_{0}u_{0}^{*})||_{L^{q}(\Omega)}$$

$$\leq ||(A_{\varepsilon} + V_{0})^{-1} - E_{\varepsilon}(A_{0} + V_{0})^{-1}M_{\varepsilon})(f(u_{\varepsilon}^{*}) + V_{0}u_{\varepsilon}^{*})||_{L^{q}(\Omega)}$$

$$+ ||E_{\varepsilon}(A_{0} + V_{0})^{-1}M_{\varepsilon}[f(u_{\varepsilon}^{*}) - V_{0}u_{\varepsilon}^{*} - E_{\varepsilon}(f(u_{0}^{*}) + V_{0}M_{\varepsilon}E_{\varepsilon}u_{0}^{*})]||_{L^{q}(\Omega)}$$

$$\leq C \varepsilon^{N/q}||f(u_{\varepsilon}^{*}) + V_{0}u_{\varepsilon}^{*}||_{L^{q}(\Omega_{\varepsilon})}$$

$$+ ||E_{\varepsilon}(A_{0} + V_{0})^{-1}M_{\varepsilon}[f(u_{\varepsilon}^{*}) - E_{\varepsilon}f(u_{0}^{*}) - V_{0}(u_{\varepsilon}^{*} - E_{\varepsilon}u_{0}^{*})]||_{L^{q}(\Omega)}$$

$$\leq C \varepsilon^{N/q} + ||E_{\varepsilon}(A_{0} + V_{0})^{-1}M_{\varepsilon}z_{\varepsilon}||_{L^{q}(\Omega)}.$$

where $z_{\varepsilon} = f(u_{\varepsilon}^*) - f(u_0^*) + V_0(u_{\varepsilon}^* - u_0^*)$ and we have used Proposition 3.7, the boundedness of f' and that u_{ε}^* is also bounded in the sup norm uniformly in ε .

We have

$$|z_{\varepsilon}(x)| = |f(u_{\varepsilon}^*(x)) - f(u_0^*(x)) + f'(E_{\varepsilon}u_0^*(x))(u_{\varepsilon}^*(x) - E_{\varepsilon}u_0^*(x))|$$

$$\leq |[f'(\chi_{\varepsilon}^*(x)) - f'(E_{\varepsilon}u_0^*(x))](u_{\varepsilon}^*(x) - E_{\varepsilon}u_0^*(x))|$$

where $\chi_{\varepsilon}^*(x) = \theta(x)u_{\varepsilon}^*(x) + (1 - \theta(x))E_{\varepsilon}u_0^*(x)$ and $0 \leqslant \theta(x) \leqslant 1$, $x \in \Omega_{\varepsilon}$. Using that $|f'(\cdot)| \leqslant C$ we have,

$$||z_{\varepsilon}||_{L^{r}(\Omega)} \leqslant C||u_{\varepsilon}^{*} - E_{\varepsilon}u_{0}^{*}||_{L^{r}(\Omega)}, \quad \forall 1 \leqslant r \leqslant +\infty.$$

Also,

$$||z_{\varepsilon}||_{L^{r}(\Omega)} \leqslant ||f'(\chi_{\varepsilon}^{*}(x)) - f'(E_{\varepsilon}u_{0}^{*}(x))||_{L^{s}(\Omega)}||u_{\varepsilon}^{*} - E_{\varepsilon}u_{0}^{*}||_{L^{t}(\Omega)}, \quad \frac{1}{r} = \frac{1}{\varepsilon} + \frac{1}{t}$$

But

$$||f'(\chi_{\varepsilon}^{*}(x)) - f'(E_{\varepsilon}u_{0}^{*}(x))||_{L^{\infty}(\Omega)} \leq C ||f'(\chi_{\varepsilon}^{*}(x)) - f'(E_{\varepsilon}u_{0}^{*}(x))||_{L^{1}(\Omega)} \leq C||\chi_{\varepsilon}^{*}(x)) - E_{\varepsilon}u_{0}^{*}(x)||_{L^{1}(\Omega)} \leq C||u_{\varepsilon}^{*} - E_{\varepsilon}u_{0}^{*}||_{L^{1}(\Omega)}.$$

Hence, using interpolation $||f'(\chi_{\varepsilon}^*(x)) - f'(E_{\varepsilon}u_0^*(x))||_{L^s(\Omega)} \leqslant C||u_{\varepsilon}^* - E_{\varepsilon}u_0^*||_{L^1(\Omega)}^{1/s}$. So

$$||z_{\varepsilon}||_{L^{r}(\Omega)} \leqslant C ||u_{\varepsilon}^{*} - E_{\varepsilon}u_{0}^{*}||_{L^{1}(\Omega)}^{1/s} ||u_{\varepsilon}^{*} - E_{\varepsilon}u_{0}^{*}||_{L^{t}(\Omega)} \leqslant C ||u_{\varepsilon}^{*} - E_{\varepsilon}u_{0}^{*}||_{L^{t}(\Omega)}^{1+\frac{1}{s}}$$

But if we define $w_{\varepsilon} = E_{\varepsilon}(A_0 + B)^{-1}M_{\varepsilon}z_{\varepsilon}$, we know from (3.16) that

$$||w_{\varepsilon}||_{L^{q}(\Omega)} \leq C||z_{\varepsilon}||_{L^{r}(\Omega)}$$
 for some $r < q$.

Hence we can choose $\frac{1}{r} = \frac{1}{s} + \frac{1}{q}$ $(t = q, \frac{1}{s} = \frac{1}{r} - \frac{1}{q} > 0)$. So

$$||E_{\varepsilon}(A_0+B)^{-1}M_{\varepsilon}z_{\varepsilon}||_{L^q(\Omega)} \leqslant C||z_{\varepsilon}||_{L^r(\Omega)} \leqslant C||u_{\varepsilon}^*-u_0^*||_{L^q(\Omega)}^{1+\frac{1}{r}-\frac{1}{q}}.$$

Hence

$$||u_{\varepsilon}^* - u_0^*||_{L^q(\Omega)} \leqslant C\varepsilon^{N/q} + C||u_{\varepsilon}^* - u_0^*||_{L^q(\Omega)}^{1 + \frac{1}{r} - \frac{1}{q}}.$$

Since we know that $||u_{\varepsilon}^* - u_0^*||_{L^q(\Omega)} \to 0$ (since $||u_{\varepsilon}^* - u_0^*||_{U_{\varepsilon}^p} \to 0$ as $\varepsilon \to 0$) then $||u_{\varepsilon}^* - u_0^*||_{L^q(\Omega)} \leqslant C \varepsilon^{N/q}$, which shows the first statement of the lemma. For the second one, we just realize that

$$||u_{\varepsilon}^* - u_0^*||_{L^q(\Omega)} + ||u_{\varepsilon}^* - u_0^*||_{L^q(R_{\varepsilon})} \leqslant C \varepsilon^{N/q} + C \mathbf{o}(\varepsilon^{\frac{N-1}{q}}) = \mathbf{o}(\varepsilon^{\frac{N-1}{q}}).$$

That is,

$$\varepsilon^{-\frac{N-1}{q}} \|u_{\varepsilon}^* - u_0^*\|_{L^q(\Omega_{\varepsilon})} \to 0 \text{ as } \varepsilon \to 0.$$

Corollary 3.16. In the conditions of Proposition 3.15, if we denote by $V_{\varepsilon} = -f'(u_{\varepsilon}^*)$, $V_0 = -f'(u_0^*)$ and $a = \sup\{|f'(s)|; s \in \mathbb{R}\}$, then hypothesis (\mathbf{H}) from Subsection 3.2 is satisfied. Hence, all the results of that Subsection can be applied to the case where the potentials are given by $V_{\varepsilon} = -f'(u_{\varepsilon}^*)$ and $V_0 = -f'(u_0^*)$.

Proof: Since

$$\|V_{\varepsilon} - E_{\varepsilon}V_0\|_{L^q(\Omega_{\varepsilon})} = \|f'(u_{\varepsilon}^*) - E_{\varepsilon}f'(u_0^*)\|_{L^q(\Omega_{\varepsilon})} \leqslant \|f''\|_{L^{\infty}(\mathbb{R})} \|u_{\varepsilon}^* - E_{\varepsilon}u_0^*\|_{L^q(\Omega_{\varepsilon})} = \mathbf{o}(\varepsilon^{\frac{N-1}{q}}),$$
 the result follows.

4. Convergence of Linear Semigroups

In this section we analyze the convergence properties of the linear semigroups generated by the operators $A_{\varepsilon} + V_{\varepsilon}$, $A_0 + V_0$ where the potentials V_{ε} , V_0 satisfy hypothesis (**H**) from Subsection 3.2. Later on we will be interested in applying the results from this section to the semigroups generated by A_{ε} , A_0 and also by $A_{\varepsilon} - f'(u_{\varepsilon}^*)$ and $A_0 - f'(u_0^*)$, where u_{ε}^* , u_0^* are hyperbolic equilibria of the perturbed and limit problem respectively.

As in Section 3.2, let $W_{\varepsilon} = V_{\varepsilon} + a > 00$, $W_0 = V_0 + a > 0$, (see hypothesis (**H**)) and $\Lambda_{\varepsilon} = A_{\varepsilon} + W_{\varepsilon}$, $\Lambda_0 = A_0 + W_0$.

As we have already seen in [4], the operators $-A_0$, $-(A_0 + V_0)$ and $-\Lambda_0$ do not generate strongly continuous semigroups in U_0^p . Nonetheless they generate certain singular semigroups as we briefly recall.

Let $\Sigma_{\theta} = \{\lambda \in \mathbb{C} : |\arg(\lambda)| \leq \pi - \theta\}$, $0 < \theta < \frac{\pi}{2}$ and let Γ be the boundary of Σ_{θ} oriented such that the imaginary part grows as λ runs in Γ . Notice that the semigroups generated by $-\Lambda_0$ and by $-(A_0 + V_0)$ are related by a multiplicative factor of the form e^{at} .

Proceeding as in [4] we define

$$e^{-\Lambda_0 t} = \frac{1}{2\pi i} \int_{\Gamma} e^{\lambda t} \left(\lambda + \Lambda_0\right)^{-1} d\lambda, \quad t > 0. \tag{4.1}$$

Then, $e^{-\Lambda_0 t}$ satisfies the semigroup properties but strong continuity fails at t=0 for data which are not sufficiently smooth. Nonetheless, several of the properties of analytic semigroup will still hold for sufficiently regular data. We say that $\{e^{-\Lambda_0 t}: t \geq 0\}$ is the semigroup generated by $-\Lambda_0$ and do not make any allusion to continuity. We refer to [4] for a detailed study of the semigroup generated by $-\Lambda_0$.

In what follows we recall some simple properties of the semigroup $\{e^{-\Lambda_0 t}: t \ge 0\}$ that we will employed later in this paper.

The next result investigates the singularity of $\{E_{\varepsilon}e^{-\Lambda_0t}M_{\varepsilon}: t>0\}$ at t=0 in $\mathcal{L}(U_{\varepsilon}^p)$. Its proof is a consequence of Proposition 3.12 and (4.1).

Lemma 4.1. For any $p \geqslant q > \frac{N}{2}$ and for $0 < \alpha < 1 - \frac{N}{2q} - \frac{1}{2}(\frac{1}{q} - \frac{1}{p}) < 1$, there is a constant C, independent of ε , such that

$$||E_{\varepsilon}e^{-\Lambda_0 t}M_{\varepsilon}u||_{U_{\varepsilon}^p} \leqslant Ct^{\alpha-1}||u||_{U_{\varepsilon}^q}, \quad t > 0, \quad u \in U_{\varepsilon}^q,$$

$$\tag{4.2}$$

and

$$||E_{\varepsilon}e^{-\Lambda_0 t}M_{\varepsilon}u||_{U_{\varepsilon}^p} \leqslant C||u||_{U_{\varepsilon}^{\infty}}, \quad t > 0, \quad u \in U_{\varepsilon}^{\infty}.$$

$$(4.3)$$

From Lemma 3.8 it follows that, $-\Lambda_{\varepsilon}$ generates an analytic semigroup $\{e^{-\Lambda_{\varepsilon}t}: t \geqslant 0\}$ in U^p_{ε} given by

$$e^{-\Lambda_{\varepsilon}t} = \frac{1}{2\pi i} \int_{\Gamma} e^{\lambda t} (\lambda + \Lambda_{\varepsilon})^{-1} d\lambda, \quad t > 0,$$
(4.4)

where $\Gamma \subset \rho(-\Lambda_{\varepsilon})$ is the boundary of Σ_{θ} oriented such that the imaginary part grows as λ runs in Γ . Note that Γ is independent of ε . It follows from (3.22), (3.23) and (4.4) that the following estimates hold

$$\|e^{-\Lambda_{\varepsilon}t}w\|_{U_{\varepsilon}^{p}} \leqslant C\varepsilon^{\frac{-N+1}{p}}\|w\|_{U_{\varepsilon}^{p}}, \quad t \geqslant 0, \ w \in U_{\varepsilon}^{p},$$
 (4.5)

$$||e^{-\Lambda_{\varepsilon}t}w||_{L^{p}(\Omega_{\varepsilon})} \leqslant C||w||_{L^{p}(\Omega_{\varepsilon})}, \quad t \geqslant 0, \ w \in L^{p}(\Omega_{\varepsilon}),$$
 (4.6)

and

$$||e^{-\Lambda_{\varepsilon}t}w||_{U_{\varepsilon}^{p}} \leqslant C||w||_{U_{\varepsilon}^{\infty}}, \quad t \geqslant 0, \ w \in U_{\varepsilon}^{\infty},$$
 (4.7)

for some constant C > 0 that does not depend on ε . That is, the linear semigroup $e^{\Lambda_{\varepsilon}t}$ is bounded in $\mathcal{L}(L^p(\Omega_{\varepsilon}))$ uniformly with respect to ε .

We analyze now the convergence properties of the semigroups. To accomplish this task we will use extensively the resolvent estimates of the previous section applied to the integral expression of the semigroup.

Proposition 4.2. There are $\gamma > 0$, $\beta \in \mathbb{R}$, p, q > N and function $\rho : [0, 1] \to [0, \infty)$ with $\rho(\varepsilon) \xrightarrow{\varepsilon \to 0} 0$ such that

$$||e^{(A_{\varepsilon}+V_{\varepsilon})t} - E_{\varepsilon}e^{(A_0+V_0)t}M_{\varepsilon}||_{\mathcal{L}(U_{\varepsilon}^p,U_{\varepsilon}^q)} \leqslant C e^{\beta t}t^{-\gamma}\rho(\varepsilon), \quad t > 0.$$
(4.8)

Proof: Observe first that $e^{-(A_{\varepsilon}+V_{\varepsilon})t} - E_{\varepsilon}e^{-(A_0+V_0)t}M_{\varepsilon} = e^{at}(e^{-\Lambda_{\varepsilon}t} - E_{\varepsilon}e^{-\Lambda_0t}M_{\varepsilon})$, so that it is sufficient to prove an estimate of the type (4.8) for the difference $e^{-\Lambda_{\varepsilon}t} - E_{\varepsilon}e^{-\Lambda_0t}M_{\varepsilon}$. Since,

$$e^{-\Lambda_{\varepsilon}t} - E_{\varepsilon}e^{-\Lambda_{0}t}M_{\varepsilon} = \frac{1}{2\pi i} \int_{\Gamma} ((\lambda + \Lambda_{\varepsilon})^{-1} - E_{\varepsilon}(\lambda + \Lambda_{0})^{-1}M_{\varepsilon})e^{\lambda t}d\lambda, \tag{4.9}$$

it follows from Proposition 3.12 that

$$\varepsilon^{-\frac{N-1}{q}} \|e^{-\Lambda_{\varepsilon}t} - E_{\varepsilon}e^{-\Lambda_{0}t}M_{\varepsilon}\|_{\mathcal{L}(U_{\varepsilon}^{p},L^{q}(\Omega_{\varepsilon}))} \leqslant \frac{C}{2\pi} \left| \int_{\Gamma} (1+|\lambda|^{1-\alpha})|e^{\lambda t}|d\lambda \right| \eta(\varepsilon)$$
$$\leqslant C t^{-(2-\alpha)}\eta(\varepsilon)$$

and consequently

$$||e^{-\Lambda_{\varepsilon}t} - E_{\varepsilon}e^{-\Lambda_{0}t}M_{\varepsilon}||_{\mathcal{L}(U_{\varepsilon}^{p},U_{\varepsilon}^{q})} \leqslant C t^{-(2-\alpha)}\eta(\varepsilon).$$

On the other hand, by comparison (maximum principle) we have

$$\|e^{-\Lambda_{\varepsilon}t} - E_{\varepsilon}e^{-\Lambda_{0}t}M_{\varepsilon}\|_{\mathcal{L}(U_{\varepsilon}^{\infty})} \leq \|e^{-\Lambda_{\varepsilon}t}\|_{\mathcal{L}(U_{\varepsilon}^{\infty})} + \|E_{\varepsilon}e^{-\Lambda_{0}t}M_{\varepsilon}\|_{\mathcal{L}(U_{\varepsilon}^{\infty})} \leq C.$$

Noting that $\|\cdot\|_{U^q_\varepsilon} \leq c \|\cdot\|_{U^\infty_\varepsilon}$ for some c>0 independent of ε , it follows that

$$||e^{-\Lambda_{\varepsilon}t} - E_{\varepsilon}e^{-\Lambda_{0}t}M_{\varepsilon}||_{\mathcal{L}(U_{\varepsilon}^{\infty},U_{\varepsilon}^{q})} \leqslant C.$$

By interpolation (see [8, Theorem 6.27])

$$\|e^{-\Lambda_{\varepsilon}} - E_{\varepsilon}e^{-\Lambda_{0}t}M_{\varepsilon}\|_{\mathcal{L}(U_{\varepsilon}^{\bar{p}},U_{\varepsilon}^{q})} \leqslant C t^{-\theta(2-\alpha)}\eta^{\theta}(\varepsilon).$$

where $p \leq \bar{p} < \infty$ and $0 \leq \theta \leq 1$. Taking θ small we can make $\theta(2 - \alpha) < 1$. That is

$$\|e^{-\Lambda_{\varepsilon}t} - E_{\varepsilon}e^{-\Lambda_{0}t}M_{\varepsilon}\|_{\mathcal{L}(U_{\varepsilon}^{\bar{p}}, U_{\varepsilon}^{q})} \leqslant C t^{-\gamma}\eta(\varepsilon)^{\theta}, \ \gamma < 1.$$

Hence, if we define $\rho(\varepsilon) = \eta(\varepsilon)^{\theta}$, we have

$$\|e^{(A_{\varepsilon}+V_{\varepsilon})t} - E_{\varepsilon}e^{(A_{0}+V_{0})t}M_{\varepsilon}\|_{\mathcal{L}(U_{\varepsilon}^{\bar{p}},U_{\varepsilon}^{q})} = e^{at}\|e^{-\Lambda_{\varepsilon}t} - E_{\varepsilon}e^{-\Lambda_{0}t}M_{\varepsilon}\|_{\mathcal{L}(U_{\varepsilon}^{\bar{p}},U_{\varepsilon}^{q})} \leqslant C e^{at}t^{-\gamma}\rho(\varepsilon)$$

which shows the result with $\rho(\varepsilon) = \eta(\varepsilon)^{\theta}$ and $\beta = a$.

Let us consider now a real number b with the property that there exists a $\delta > 0$, small, such that $[b-\delta,b+\delta] \cap \sigma(-(A_0+V_0)) = \emptyset$. That is, the spectrum of the operator $-(A_0+V_0)$, which is all real, is divided in two parts, σ_0^+ which is above $b+\delta$ and it is a finite set and σ_0^- which is below $b-\delta$ and it is an infinite set (a sequence that goes to $-\infty$). From the continuity properties of the spectrum, (see [3]) we have that for ε small enough $[b-\delta,b+\delta] \cap \sigma(-(A_\varepsilon+V_\varepsilon)) = \emptyset$ and the spectra of $-(A_\varepsilon+V_\varepsilon)$, which is also real, is divided in two parts σ_ε^+ , above $b+\delta$ and σ_ε^- , below $b-\delta$. Moreover, we can choose a fixed closed curve $\Gamma_b^+ \subset \{z \in \mathbb{C} : \operatorname{Re}(z) \geqslant b+\delta\}$ which encloses σ_ε^+ for all $0 \leqslant \varepsilon \leqslant \varepsilon_0$ for some ε_0 small. Moreover, we denote by $\Gamma_b^- = \{z \in \mathbb{C} : \operatorname{arg}(z-(b-\delta)) = \pi-\theta\}$ for some $0 < \theta < \pi/2$.

We decompose U_{ε}^{p} using the projection

$$Q_{\varepsilon}^{+} = Q(\sigma_{\varepsilon}^{+}) = \frac{1}{2\pi i} \int_{\Gamma_{h}^{+}} (\lambda + A_{\varepsilon} + V_{\varepsilon})^{-1} d\lambda. \tag{4.10}$$

Proposition 4.3. For p, q > N large enough, we have that there are constants C > 0, $\gamma < 1$, independent of ε and a function $\rho(\varepsilon)$, with $\rho(\varepsilon) \to 0$ as $\varepsilon \to 0$, such that for t > 0

$$||e^{-(A_{\varepsilon}+V_{\varepsilon})t}(I-Q(\sigma_{\varepsilon}^{+}))-E_{\varepsilon}e^{-(A_{0}+V_{0})t}(I-Q(\sigma_{0}^{+}))M_{\varepsilon}||_{\mathcal{L}(U_{\varepsilon}^{p},U_{\varepsilon}^{q})} \leqslant C e^{bt}t^{-\gamma}\rho(\varepsilon)$$
(4.11)

$$||E_{\varepsilon}e^{-(A_0+V_0)t}(I-Q(\sigma_0^+))M_{\varepsilon}||_{\mathcal{L}(U_{\varepsilon}^p,U_{\varepsilon}^q)} \leqslant C e^{bt}t^{-\gamma}$$
(4.12)

$$||e^{-(A_{\varepsilon}+V_{\varepsilon})t}(I-Q(\sigma_{\varepsilon}^{+}))||_{\mathcal{L}(U_{\varepsilon}^{p},U_{\varepsilon}^{q})} \leqslant C e^{bt} t^{-\gamma}$$

$$\tag{4.13}$$

Proof: We have

$$e^{-(A_0+V_0)t}(I-Q(\sigma_0^+)) = \frac{1}{2\pi i} \int_{\Gamma_*^-} (\lambda + A_0 + V_0(x))^{-1} e^{\lambda t} d\lambda.$$

Plugging norms and using estimate (3.5) we get

$$||e^{-(A_0+V_0)t}(I-Q(\sigma_0^+))||_{\mathcal{L}(U_0^p,U_0^q)} \leqslant \left|\frac{1}{2\pi} \int_{\Gamma_b^-} \frac{|e^{\lambda t}|}{|1+|\lambda|^{1-\alpha}} d\lambda\right|$$

and elementary integration shows

$$||e^{(A_0+V_0)t}(I-Q(\sigma_0^+))||_{\mathcal{L}(U_0^p,U_0^q)} \leqslant Ce^{bt}t^{-\alpha}$$
(4.14)

which shows (4.12) with $\gamma = \alpha$.

In a similar way,

$$e^{-(A_{\varepsilon}+V_{\varepsilon})t}(I-Q(\sigma_{\varepsilon}^{+})) - E_{\varepsilon}e^{-(A_{0}+V_{0})t}(I-Q(\sigma_{0}^{+}))M_{\varepsilon} = \frac{1}{2\pi i} \int_{\Gamma_{h}^{-}} ((\lambda+A_{\varepsilon}+V_{\varepsilon}(x))^{-1} - E_{\varepsilon}(\lambda+A_{0}+V_{0}(x))^{-1}M_{\varepsilon})e^{\lambda t}d\lambda.$$

So

$$\begin{aligned} \|e^{-(A_{\varepsilon}+V_{\varepsilon})t}(I-Q(\sigma_{\varepsilon}^{+})) - E_{\varepsilon}e^{-(A_{0}+V_{0})t}(I-Q(\sigma_{0}^{+}))M_{\varepsilon}\|_{\mathcal{L}(U_{\varepsilon}^{p},U_{\varepsilon}^{q})} &\leq \\ \frac{1}{2\pi} \left| \int\limits_{\Gamma_{b}^{-}} |e^{\lambda t}| \|(\lambda+A_{\varepsilon}+V_{\varepsilon}(x))^{-1} - E_{\varepsilon}(\lambda+A_{0}+V_{0}(x))^{-1}M_{\varepsilon})\|_{\mathcal{L}(U_{\varepsilon}^{p},U_{\varepsilon}^{q})} d\lambda \right| \\ &\leq \frac{1}{2\pi} \int\limits_{\Gamma_{b}^{-}} |e^{\lambda t}| (1+|\lambda|^{1-\alpha}) d\lambda \eta(\varepsilon) d\lambda \\ &\leq \frac{C}{2\pi} e^{bt} t^{-(2-\alpha)} \eta(\varepsilon), \end{aligned}$$

where we have applied Proposition 3.12. Therefore,

$$\|e^{-(A_{\varepsilon}+V_{\varepsilon})t}(I-Q(\sigma_{\varepsilon}^{+}))-E_{\varepsilon}e^{-(A_{0}+V_{0})t}(I-Q(\sigma_{0}^{+}))M_{\varepsilon}\|_{\mathcal{L}(U_{\varepsilon}^{p},U_{\varepsilon}^{q})} \leqslant C e^{bt}t^{-(2-\alpha)}\eta(\varepsilon). \tag{4.15}$$

This estimate does not show yet the proposition since the exponent $2 - \alpha > 1$. We will do an interpolation argument to conclude with the correct estimate. For this, let us see now that $Q(\sigma_{\varepsilon}^+): U_{\varepsilon}^p \to U_{\varepsilon}^p$ satisfies $\|Q(\sigma_{\varepsilon}^+)\|_{\mathcal{L}(U_{\varepsilon}^p,U_{\varepsilon}^p)} \leqslant C$ independent of ε . To see this, just observe that

$$Q(\sigma_{\varepsilon}^{+}) = \frac{1}{2\pi i} \int_{\Gamma_{h}^{+}} (\lambda + A_{\varepsilon} + V_{\varepsilon})^{-1} d\lambda.$$

Applying now the estimate of Propostion 3.12, we obtain that

$$\|(\lambda + A_{\varepsilon} + V_{\varepsilon})^{-1}\|_{\mathcal{L}(U_{\varepsilon}^{p}, U_{\varepsilon}^{p})} \le C$$

for $\lambda \in \Gamma_b^-$ and with C independent of ε . From this last expression and using the boundedness of Γ_b^- we get $\|Q(\sigma_\varepsilon^+)\|_{\mathcal{L}(U_\varepsilon^p,U_\varepsilon^p)} \leqslant C$, for all $0 \leqslant \varepsilon \leqslant 1$.

Moreover, for the limit semigroup and for $0 < t \le 1$, we obtain from (4.14)

$$||E_{\varepsilon}e^{-(A_0+V_0)t}(I-Q(\sigma_0^+))M_{\varepsilon}||_{\mathcal{L}(U_{\varepsilon}^p,U_{\varepsilon}^q)} \leqslant Ct^{-\alpha}.$$

Hence for $0 < t \le 1$, we get that

$$\|e^{-(A_{\varepsilon}+V_{\varepsilon})t}(I-Q(\sigma_{\varepsilon}^{+}))\|_{\mathcal{L}(U_{\varepsilon}^{p},U_{\varepsilon}^{q})} \leqslant \|e^{-(A_{\varepsilon}+V_{\varepsilon})t}\|_{\mathcal{L}(U_{\varepsilon}^{p},U_{\varepsilon}^{q})}(1+\|Q(\sigma_{\varepsilon}^{+})\|_{\mathcal{L}(U_{\varepsilon}^{p},U_{\varepsilon}^{p})})$$

$$C\left(\|e^{-(A_{\varepsilon}+V_{\varepsilon})t}-E_{\varepsilon}e^{-(A_{0}+V_{0})t}M_{\varepsilon}\|_{\mathcal{L}(U_{\varepsilon}^{p},U_{\varepsilon}^{q})}+\|E_{\varepsilon}e^{-(A_{0}+V_{0})t}M_{\varepsilon}\|_{\mathcal{L}(U_{\varepsilon}^{p},U_{\varepsilon}^{q})}\right) \leqslant C(t^{-\gamma}+t^{-\alpha+1})$$

where we are using the bounds given by Proposition 4.2

Hence, for $0 < t \le 1$

$$\|e^{-(A_{\varepsilon}+V_{\varepsilon})t}(I-Q(\sigma_{\varepsilon}^{+}))-E_{\varepsilon}e^{-(A_{0}+V_{0})t}(I-Q(\sigma_{0}^{+}))M_{\varepsilon}\|_{\mathcal{L}(U_{\varepsilon}^{p},U_{\varepsilon}^{q})} \leqslant Ct^{-\bar{\gamma}}$$

$$\tag{4.16}$$

where $\bar{\gamma} = \max\{\gamma, 1 - \alpha\}.$

Interpolating (4.15) and (4.16) we obtain, for $0 < t \le 1$,

$$||e^{-(A_{\varepsilon}+V_{\varepsilon})t}(I-Q(\sigma_{\varepsilon}^{+})) - E_{\varepsilon}e^{-(A_{0}+V_{0})t}(I-Q(\sigma_{0}^{+}))M_{\varepsilon}||_{\mathcal{L}(U_{\varepsilon}^{p},U_{\varepsilon}^{q})} \leqslant (4.17)$$

$$(Ct^{-(2-\alpha)}\eta(\varepsilon))^{\theta}(Ct^{-\bar{\gamma}})^{1-\theta} \leqslant Ct^{-(2-\alpha)\theta-(1-\theta)\bar{\gamma}}\eta(\varepsilon)^{\theta}$$
(4.18)

where we have used that $e^{bt} \leqslant C$ for $0 \leqslant t \leqslant 1$. Choosing $\theta > 0$ small enough so that $(2-\alpha)\theta + (1-\theta)\bar{\gamma} < 1$, we obtain the estimate for $0 < t \leqslant 1$.

Now for $t \ge 1$, from (4.15) we get

$$\|e^{-(A_{\varepsilon}+V_{\varepsilon})t}(I-Q(\sigma_{\varepsilon}^{+}))-E_{\varepsilon}e^{-(A_{0}+V_{0})t}(I-Q(\sigma_{0}^{+}))M_{\varepsilon}\|_{\mathcal{L}(U_{\varepsilon}^{p},U_{\varepsilon}^{q})}\leqslant Ce^{bt}\eta(\varepsilon)$$

Putting together both estimates, we prove (4.11). To prove (4.13) we just use (4.11) and (4.12). This concludes the proof of the proposition.

We also have

Corollary 4.4. For the case $V_{\varepsilon} = V_0 \equiv 0$ and with $b \in (-1,0)$ a fixed number, we have that $Q(\sigma_{\varepsilon}^+) \equiv 0$ for ε small enough and we have

$$||e^{-A_{\varepsilon}t} - E_{\varepsilon}e^{-A_{0}t}M_{\varepsilon}||_{\mathcal{L}(U_{\varepsilon}^{p}, U_{\varepsilon}^{q})} \leqslant C e^{bt}t^{-\gamma}\rho(\varepsilon).$$

Remark 4.5. Observe that we can consider the case where $V_0 = -f'(u_0^*)$, $V_{\varepsilon} = -f'(u_{\varepsilon}^*)$ with u_0^* and u_{ε}^* hyperbolic equilibria satisfying u_{ε}^* converging to u_0^* (see [3]). In this case, we can always apply Proposition 4.3 with b < 0, a number dividing the spectrum among the stable part, that is with negative real part, and the unstable spectrum, that is with positive real part.

Let us conclude the section with the following useful unifom estimates of the semigroup on the linear unstable manifold

Proposition 4.6. There are constants $C \geqslant 1$ and $\beta > 0$ such that

$$||e^{-(A_{\varepsilon}+V_{\varepsilon})t}Q_{\varepsilon}^{+}||_{\mathcal{L}(U_{\varepsilon}^{q},U_{\varepsilon}^{p})} \leqslant Ce^{\beta t}, \ t \leqslant 0$$

$$\tag{4.19}$$

Proof: Observe that

$$e^{-(A_{\varepsilon}+V_{\varepsilon})t}Q_{\varepsilon}^{+} = \int_{\Gamma^{+}} e^{\lambda t} (\lambda + A_{\varepsilon} + V_{\varepsilon})^{-1} d\lambda.$$

Using (3.36) and noticing that the curve Γ^+ is bounded, we have

$$\|e^{-(A_{\varepsilon}+V_{\varepsilon})t}Q_{\varepsilon}^{+}\|_{\mathcal{L}(U_{\varepsilon}^{q},U_{\varepsilon}^{p})} \leqslant C \left| \int_{\Gamma^{+}} |e^{\lambda t}| d\lambda \right| \leqslant Ce^{\beta t}$$

which shows the result.

5. Continity of nonlinear semigroups and upper semicontinuity of attractors

Now that we have obtained in the previous section the continuity of linear semigroups we proceed to obtain the continuity of nonlinear semigroups using the Variation of Constants Formula. After we obtain the continuity of nonlinear semigroups we will proceed to obtain the upper semicontinuity of the family of attractors $\{A_{\varepsilon} : \varepsilon \in [0, 1]\}$.

To this end we will follow the ideas in [1] that relate the continuity of the linear semigroups with the continuity of the nonlinear semigroups for dissipative parabolic equations by using the variation of constants formula. This in turn will imply the upper semicontinuity of the attractors and the stationary states.

For $\varepsilon \in [0,1]$, let $\{T_{\varepsilon}(t): t \geq 0\}$ be the semigroups defined in U_{ε}^p by the variation of constants formula

$$T_{\varepsilon}(t, u_{\varepsilon}) = e^{-A_{\varepsilon}t}u_{\varepsilon} + \int_{0}^{t} e^{-A_{\varepsilon}(t-s)} f_{\varepsilon}(T_{\varepsilon}(s, u_{\varepsilon}))ds.$$
 (5.1)

If $\mathcal{E}_{\varepsilon}$ denotes the set of stationary states (2.6), $\varepsilon \in [0, \varepsilon_0]$, it has been obtained in [3, Section 5] that, $\{\mathcal{E}_{\varepsilon} : \varepsilon \in [0, \varepsilon_0]\}$ is upper semicontinuous at $\varepsilon = 0$ in U_{ε}^p ; that is,

$$\sup_{u_{\varepsilon}^* \in \mathcal{E}_{\varepsilon}} \left[\inf_{u_0^* \in \mathcal{E}_0} \{ \|u_{\varepsilon}^* - E_{\varepsilon} u_0^*\|_{U_{\varepsilon}^p} \} \right] \to 0, \quad \text{as } \varepsilon \to 0$$
 (5.2)

We are now in position to prove the following result

Proposition 5.1. There exists a $0 \le \gamma < 1$ and a function $c(\varepsilon)$ with $c(\varepsilon) \xrightarrow{\varepsilon \to 0} 0$ such that, for each $\tau > 0$ we have

$$||T_{\varepsilon}(t, u_{\varepsilon}) - E_{\varepsilon}T_{0}(t, M_{\varepsilon}u_{\varepsilon})||_{U_{\varepsilon}^{p}} \leqslant M(\tau)c(\varepsilon)t^{-\gamma}, \quad t \in (0, \tau], \quad u_{\varepsilon} \in \mathcal{A}_{\varepsilon}, \quad \varepsilon \in (0, \varepsilon_{0}]. \quad (5.3)$$

Moreover, the family of attractors $\{A_{\varepsilon} : \varepsilon \in [0, \varepsilon_0]\}$ is upper semicontinuous at $\varepsilon = 0$ in U_{ε}^p , in the sense that

$$\sup_{u_{\varepsilon} \in \mathcal{A}_{\varepsilon}} \left[\inf_{u_{0} \in \mathcal{A}_{0}} \{ \|u_{\varepsilon} - E_{\varepsilon} u_{0}\|_{U_{\varepsilon}^{p}} \} \right] \to 0, \quad as \ \varepsilon \to 0$$
 (5.4)

Proof: To prove this result we follow [1, 7]. Notice that the nonlinear semigroups $T_{\varepsilon}(t)$ are given by (5.1). Hence, estimating $T_{\varepsilon}(t, u_{\varepsilon}) - E_{\varepsilon}T_0(t, M_{\varepsilon}u_{\varepsilon})$ and with some elementary computations we obtain

$$\begin{split} \|T_{\varepsilon}(t,u_{\varepsilon}) - E_{\varepsilon}T_{0}(t,M_{\varepsilon}u_{\varepsilon})\|_{U_{\varepsilon}^{p}} &\leq \|e^{-A_{\varepsilon}t}u_{\varepsilon} - E_{\varepsilon}e^{-A_{0}t}M_{\varepsilon}u_{\varepsilon}\|_{U_{\varepsilon}^{p}} \\ &+ \int_{0}^{t} \|\left(e^{-A_{\varepsilon}t} - E_{\varepsilon}e^{-A_{0}t}M_{\varepsilon}\right)f_{\varepsilon}(T_{\varepsilon}(s,u_{\varepsilon}))\|_{U_{\varepsilon}^{p}} ds \\ &+ \int_{0}^{t} \|E_{\varepsilon}e^{-A_{0}t}(M_{\varepsilon}f_{\varepsilon}(T_{\varepsilon}(s,u_{\varepsilon})) - f_{0}(T_{0}(s,M_{\varepsilon}u_{\varepsilon})))\|_{U_{\varepsilon}^{p}} ds, \quad \varepsilon \in [0,\varepsilon_{0}]. \end{split}$$

Note that

$$\int_{0}^{t} \|E_{\varepsilon}e^{-A_{0}t}(M_{\varepsilon}f_{\varepsilon}(T_{\varepsilon}(s,u_{\varepsilon})) - f_{0}(T_{0}(s,M_{\varepsilon}u_{\varepsilon})))\|_{U_{\varepsilon}^{p}} ds$$

$$= \int_{0}^{t} \|E_{\varepsilon}e^{-A_{0}t}(M_{\varepsilon}f_{\varepsilon}(T_{\varepsilon}(s,u_{\varepsilon})) - M_{\varepsilon}E_{\varepsilon}f_{0}(T_{0}(s,M_{\varepsilon}u_{\varepsilon})))\|_{U_{\varepsilon}^{p}} ds$$

$$= \int_{0}^{t} \|E_{\varepsilon}e^{-A_{0}t}M_{\varepsilon}(f_{\varepsilon}(T_{\varepsilon}(s,u_{\varepsilon})) - f_{\varepsilon}(E_{\varepsilon}T_{0}(s,M_{\varepsilon}u_{\varepsilon})))\|_{U_{\varepsilon}^{p}} ds$$

where we have used that $M_{\varepsilon}E_{\varepsilon} = I$ and that $f_{\varepsilon}(E_{\varepsilon}u) = E_{\varepsilon}f_{0}(u)$. Applying now Corollary 4.4 and Lemma 4.1 we have, for $0 < t \le \tau$,

$$||T_{\varepsilon}(t, u_{\varepsilon}) - E_{\varepsilon}T_{0}(t, M_{\varepsilon}u_{\varepsilon})||_{U_{\varepsilon}^{p}} \leqslant Ce^{bt}t^{-\gamma}\rho(\varepsilon)||u_{\varepsilon}||_{U_{\varepsilon}^{p}} + C\rho(\varepsilon)\int_{0}^{t}(t-s)^{-\gamma}e^{b(t-s)}||f_{\varepsilon}(T_{\varepsilon}(s, u_{\varepsilon}))||_{U_{\varepsilon}^{p}} + C\int_{0}^{t}(t-s)^{\alpha-1}||T_{\varepsilon}(s, u_{\varepsilon}) - E_{\varepsilon}T_{0}(s, M_{\varepsilon}u_{\varepsilon})||_{U_{\varepsilon}^{p}}$$

But since we have uniform bounds in $L^{\infty}(\Omega_{\varepsilon})$ of all the attractors, the first two terms in the above inequality can be bounded by $C\rho(\varepsilon)t^{-\gamma}$. The result now follows applying the singular Gronwall's lemma (see [11]).

To show the uppersemicontinuity of the attractors $\mathcal{A}_{\varepsilon}$, we notice first that by the uniform $L^{\infty}(\Omega_{\varepsilon})$ bounds of the attractors we have

$$\bigcup_{0 \leqslant \varepsilon \leqslant \varepsilon_0} M_{\varepsilon} \mathcal{A}_{\varepsilon}$$

is a bounded set in U_0^{∞} . Hence, by the attractivity properties of \mathcal{A}_0 , for a fixed $\eta > 0$ there exists a time $\tau > 0$ such that

$$\operatorname{dist}_{U_0^p}\Big(T_0(\tau)(M_{\varepsilon}\varphi_{\varepsilon}), \mathcal{A}_0\Big) \equiv \inf_{\varphi \in \mathcal{A}_0} \|T_0(\tau)(M_{\varepsilon}\varphi_{\varepsilon}) - \varphi\|_{U_0^p} \leqslant \eta, \quad \forall \varphi_{\varepsilon} \in \mathcal{A}_{\varepsilon}, \quad 0 \leqslant \varepsilon \leqslant \varepsilon_0$$

which implies that

$$\operatorname{dist}_{U_{\varepsilon}^{p}}\left(E_{\varepsilon}T_{0}(\tau)(M_{\varepsilon}\varphi_{\varepsilon}), E_{\varepsilon}\mathcal{A}_{0}\right) \leqslant \eta, \quad \forall \varphi_{\varepsilon} \in \mathcal{A}_{\varepsilon}, \quad 0 \leqslant \varepsilon \leqslant \varepsilon_{0}$$

Using the convergence of the nonlinear semigroups (5.3) with $t = \tau$, there exists $\varepsilon_1 > 0$ such that for $0 < \varepsilon \leqslant \varepsilon_1$,

$$||T_{\varepsilon}(\tau, \varphi_{\varepsilon}) - E_{\varepsilon}T_{0}(\tau, M_{\varepsilon}\varphi_{\varepsilon})||_{U_{\varepsilon}^{p}} \leqslant \eta, \quad \forall \varphi_{\varepsilon} \in \mathcal{A}_{\varepsilon}, \quad 0 \leqslant \varepsilon \leqslant \varepsilon_{1}.$$

Hence,

$$\operatorname{dist}_{U_{\varepsilon}^{p}}\Big(T_{\varepsilon}(\tau,\varphi_{\varepsilon}), E_{\varepsilon}\mathcal{A}_{0}\Big) \leqslant \eta, \quad \forall \varphi_{\varepsilon} \in \mathcal{A}_{\varepsilon}, \quad 0 \leqslant \varepsilon \leqslant \varepsilon_{1}.$$

From the invariance A_{ε} we have that

$$\operatorname{dist}_{U_{\varepsilon}^{p}}\left(\varphi_{\varepsilon}, E_{\varepsilon} \mathcal{A}_{0}\right) \leqslant \eta, \quad \forall \varphi_{\varepsilon} \in \mathcal{A}_{\varepsilon}, \quad 0 \leqslant \varepsilon \leqslant \varepsilon_{1}$$

which implies (5.4).

Remark 5.2. Observe that Proposition 5.1 proves the upper semicontinuity part of Theorem 2.5.

6. Continuity of local unstable manifolds and of attractors

We already know that, if all equilibrium points of (2.7), which is the abstract version of (1.2), are hyperbolic then they are all isolated and there is only a finite number of them, say $\mathcal{E}_0 = \{e_0^1, \dots, e_0^m\}$. In this case, we also know that there is an $\varepsilon_0 > 0$ such that the set of equilibria of (2.6), which is the abstract version of (1.1), $\mathcal{E}_{\varepsilon} = \{e_{\varepsilon}^1, \dots, e_{\varepsilon}^m\}$ for all $0 < \varepsilon \leqslant \varepsilon_0$ and $e_{\varepsilon}^i \xrightarrow{E} e_0^i$ for $1 \leqslant i \leqslant m$ (see Theorem 2.3 of [3]). Moreover, we also know that the linear unstable manifolds associated to e_j^{ε} converge to the linear unstable manifold of e_j^{ε} , see Theorem 2.5 of [3]. For each $e_{\varepsilon}^j \in \mathcal{E}_{\varepsilon}$, $\varepsilon \in [0, 1]$, we define its unstable manifold

$$W^{u}(e^{j}_{\varepsilon}) = \{ \eta_{\varepsilon} \in U^{p}_{\varepsilon} : \text{ there is a global solution } \xi_{\varepsilon} : \mathbb{R} \to U^{p}_{\varepsilon} \text{ of}$$

$$(2.6) \text{ with } \xi_{\varepsilon}(0) = \eta_{\varepsilon} \text{ such that } \xi_{\varepsilon}(t) \stackrel{t \to -\infty}{\longrightarrow} e^{j}_{\varepsilon} \}.$$

and its δ -local unstable manifold as

$$W^{u}_{\delta}(e^{j}_{\varepsilon}) = \{ \eta_{\varepsilon} \in B(e^{j}_{\varepsilon}, \delta) \subset U^{p}_{\varepsilon} : \text{ there is a global solution } \xi_{\varepsilon} : \mathbb{R} \to U^{p}_{\varepsilon} \text{ of}$$

$$(2.6) \text{ with } \xi_{\varepsilon}(0) = \eta_{\varepsilon}, \ \xi_{\varepsilon}(t) \in B(e^{j}_{\varepsilon}, \delta), \ \forall \ t \leq 0, \text{ and } \xi_{\varepsilon}(t) \stackrel{t \to -\infty}{\longrightarrow} e^{j}_{\varepsilon} \}.$$

These definitions are standard and we refer to [9] for further properties of local unstable manifolds.

In this section we show that the local unstable manifolds of e_{ε}^{j} , for $j=1,\ldots,m$ fixed, behaves continuously with ε in U_{ε}^{p} .

Proposition 6.1. Assume that $e_0 \in \mathcal{E}_0$ is hyperbolic; that is, $0 \notin \sigma(A_0 - f'(e_0)I)$. By Theorem 5.8 and Example 5.9 in [3], there are $\delta > 0$ and ε_0 such that, there is a unique $e_{\varepsilon} \in \mathcal{E}_{\varepsilon}$ with $\|e_{\varepsilon} - E_{\varepsilon}e_0\|_{U_{\varepsilon}^p} < \delta$, for all $0 \leqslant \varepsilon \leqslant \varepsilon_0$. Then, there is $\delta > 0$ such that

$$\operatorname{dist}_{U_{\varepsilon}^{p}}(W_{\delta}^{u}(e_{\varepsilon}), E_{\varepsilon}W_{\delta}^{u}(e_{0})) + \operatorname{dist}_{U_{\varepsilon}^{p}}(E_{\varepsilon}W_{\delta}^{u}(e_{0}), W_{\delta}^{u}(e_{\varepsilon})) \stackrel{\varepsilon \to 0}{\longrightarrow} 0$$

that is,

$$\sup_{u_{\varepsilon} \in W^{u}_{\varepsilon}(e_{\varepsilon})} \inf_{u_{0} \in W^{u}_{\delta}(u_{0})} \|u_{\varepsilon} - E_{\varepsilon}u_{0}\|_{U^{p}_{\varepsilon}} + \sup_{u_{0} \in W^{u}_{\varepsilon}(u_{0})} \inf_{u_{\varepsilon} \in W^{u}_{\delta}(e_{\varepsilon})} \|u_{\varepsilon} - E_{\varepsilon}u_{0}\|_{U^{p}_{\varepsilon}} \to 0, \ as \ \varepsilon \to 0$$

Before proving this result, let us see how we can proceed to give a proof of our main result, Theorem 2.5.

Proof of Theorem 2.5: The upper-semicontinuity has already been proved in Proposition 5.1 from Section 5. Observe that to obtain the upper-semicontinuity of the attractors, we have used the continuity of the nonlinear semigroups, but no gradient structure of the flows have been used.

To obtain the lower-semicontinuity, we need to show that for each $\varphi_0 \in \mathcal{A}_0$ we have a sequence of $\varphi_{\varepsilon} \in \mathcal{A}_{\varepsilon}$, with the property that $\|\varphi_{\varepsilon} - E_{\varepsilon}\varphi_0\|_{U_{\varepsilon}^p} \to 0$ as $\varepsilon \to 0$. To accomplish this, we follow similar arguments as the one developed in [9], [10] or [2].

We are assuming that each equilibrium of the limiting problem \mathcal{E}_0 is hperbolic. This implies that we have a finite number of them and that the flow $T_0(t)$ has a gradient structure, see [4] and in particular, given $\varphi_0 \in \mathcal{A}_0$ it will lie in the unstable manifold of some $e_0 \in \mathcal{E}_0$. This implies that there exist an element $\phi_0 \in W^u_\delta(e_0)$ and a $\tau > 0$ such that $T_0(\tau, \phi_0) = \varphi_0$, where $\delta > 0$ is the one from Proposition 6.1. Using the continuity of the local unstable manifolds

obtained in Proposition 6.1, we have that there exists a sequence of elements $\phi_{\varepsilon} \in W^{u}_{\delta}(e_{\varepsilon})$ such that $\|\phi_{\varepsilon} - E_{\varepsilon}\phi_{0}\|_{U^{p}_{\varepsilon}} \to 0$. But, from the invariance of the attactor $\mathcal{A}_{\varepsilon}$ under the flow T_{ε} , we have $\varphi_{\varepsilon} = T_{\varepsilon}(\tau, \phi_{\varepsilon}) \in \mathcal{A}_{\varepsilon}$. Moreover,

$$\|\varphi_{\varepsilon} - E_{\varepsilon}\varphi_{0}\|_{U_{\varepsilon}^{p}} = \|T_{\varepsilon}(\tau,\phi_{\varepsilon}) - E_{\varepsilon}T_{0}(\tau,\phi_{0})\|_{U_{\varepsilon}^{p}}$$

$$\leq \|T_{\varepsilon}(\tau,\phi_{\varepsilon}) - E_{\varepsilon}T_{0}(\tau,M_{\varepsilon}\phi_{\varepsilon})\|_{U_{\varepsilon}^{p}} + \|E_{\varepsilon}T_{0}(\tau,M_{\varepsilon}\phi_{\varepsilon}) - E_{\varepsilon}T_{0}(\tau,\phi_{0})\|_{U_{\varepsilon}^{p}}$$

$$\leq M(\tau)\tau^{-\gamma}c(\varepsilon) + \|T_{0}(\tau,M_{\varepsilon}\phi_{\varepsilon}) - T_{0}(\tau,\phi_{0})\|_{U_{0}^{p}}$$

where we are using (5.3) and the fact that $||E_{\varepsilon}||_{\mathcal{L}(U_0^p,U_{\varepsilon}^p)}=1$.

The continuity of the map $T(\tau, \cdot): U_0^p \to U_0^p$, the fact that $\|\phi_0 - M_{\varepsilon}\phi_{\varepsilon}\|_{U_0^p} \to 0$ as $\varepsilon \to 0$ and that $c(\varepsilon) \to 0$, shows that $\|\varphi_{\varepsilon} - E_{\varepsilon}\varphi_0\|_{U_{\varepsilon}^p} \to 0$ as $\varepsilon \to 0$. This concludes the proof of Theorem 2.5.

Proof of Proposition 6.1: Let $\{e_{\varepsilon}\}$ with $e_{\varepsilon} \in \mathcal{E}_{\varepsilon}$, $\varepsilon \in [0,1]$, such $\|e_{\varepsilon} - E_{\varepsilon}e_0\|_{U_{\varepsilon}^p} \stackrel{\varepsilon \to 0}{\to} 0$. Rewriting (2.6) for $w_{\varepsilon} = u_{\varepsilon} - e_{\varepsilon}$ to deal with the neighborhood of e_{ε} we arrive at

$$w_t + A_{\varepsilon}w - f'_{\varepsilon}(e_{\varepsilon})w = f(w + e_{\varepsilon}) - f(e_{\varepsilon}) - f'_{\varepsilon}(e_{\varepsilon})w. \tag{6.1}$$

Let us denote by $V_0 = -f'(e_0)$, $V_{\varepsilon} = -f'(e_{\varepsilon})$. Using the hiperbolicity of e_0 , e_{ε} we consider b < 0 and define σ_{ε}^+ , $Q(\sigma_{\varepsilon}^*)$ as in (4.10), see Remark 4.5.

Decomposing (6.1) with the aid of projection $Q(\sigma_{\varepsilon}^+)$ and denoting by \tilde{A}_{ε} the restriction of $A_{\varepsilon} + V_{\varepsilon}$ to the kernel of $Q(\sigma_{\varepsilon}^+)$, by B_{ε} the restriction of $A_{\varepsilon} + V_{\varepsilon}$ to the range of $Q(\sigma_{\varepsilon}^+)$ and making $S_{\varepsilon}^{-1}v = Q(\sigma_{\varepsilon}^+)w$, $z = (I - Q(\sigma_{\varepsilon}^+))w$ we rewrite (6.1) as

$$\dot{v} + B_{\varepsilon}v = Q(\sigma_{\varepsilon}^{+})F_{\varepsilon}(S_{\varepsilon}v, z)
\dot{z} + \tilde{A}_{\varepsilon}z = (I - Q(\sigma_{\varepsilon}^{+}))F_{\varepsilon}(S_{\varepsilon}v, z),$$
(6.2)

where $F_{\varepsilon}(0,0) = 0$ and $F'_{\varepsilon}(0,0) = 0$. Proceeding as in Example 5.9 in [3] we have that, given $\rho > 0$ there is a $\delta > 0$ such that

$$||F_{\varepsilon}(S_{\varepsilon}v,z)||_{U_{\varepsilon}^{q}} < \rho, ||F_{\varepsilon}(S_{\varepsilon}v,z) - F_{\varepsilon}(S_{\varepsilon}\tilde{v},\tilde{z})||_{U_{\varepsilon}^{q}} < \rho(||v - \tilde{v}||_{\mathbb{R}^{n}} + ||z - \tilde{z}||_{U_{\varepsilon}^{p}}).$$

$$(6.3)$$

for all $(v, z) \in B_{\delta}(0, 0)$ and for all $\varepsilon \in (0, 1]$. Since we are interested only in the behavior of the solutions near (0, 0) we cut F_{ε} outside $B_{\delta}(0, 0)$ in such a way that it satisfies (6.3) globally.

Proceeding as in [2, 7] we can show that for a suitably small $\rho > 0$, there is an unstable manifold for e_{ε}

$$S^{\varepsilon} = \{(v, z) : z = \Sigma_{\varepsilon}^*(v), \ v \in \mathbb{R}^n\}$$

where $\Sigma_{\varepsilon}^* : \mathbb{R}^n \to \operatorname{Ker}(Q_{\varepsilon})$ is bounded and Lipschitz continuous. Furthermore

$$\sup_{v \in \mathbb{R}^n} \| \Sigma_{\varepsilon}^*(v) - E_{\varepsilon} \Sigma_0^*(v) \|_{U_{\varepsilon}^p} \stackrel{\varepsilon \to 0}{\to} 0.$$

Let us sketch the proof of existence of the unstable manifold as a graph and prove its continuity. Let $\Sigma_{\varepsilon} : \mathbb{R}^n \to \operatorname{Ker}(Q_{\varepsilon})$ such that

$$\|\!|\!|\!| \Sigma_{\varepsilon} \|\!|\!| := \sup_{v \in \mathbb{R}^n} \|\Sigma_{\varepsilon}(v)\|_{U_{\varepsilon}^p} \leqslant D, \ \|\Sigma_{\varepsilon}(v) - \Sigma_{\varepsilon}(\tilde{v})\|_{U_{\varepsilon}^p} \leqslant L \|v - \tilde{v}\|_{\mathbb{R}^n}. \tag{6.4}$$

If $v_{\varepsilon}(t) = \psi(t, \tau, \eta, \Sigma_{\varepsilon})$ denotes the solution of

$$\frac{dv_{\varepsilon}}{dt} + B_{\varepsilon}v_{\varepsilon} = F_{\varepsilon}(S_{\varepsilon}v_{\varepsilon}, \Sigma_{\varepsilon}(v_{\varepsilon})), \text{ for } t < \tau, \ v_{\varepsilon}(\tau) = \eta,$$

We seek for a fixed point Σ_{ε}^* of

$$\Phi(\Sigma_{\varepsilon})(\eta) = \int_{-\infty}^{\tau} e^{-\tilde{A}_{\varepsilon}(\tau - s)} (I - Q(\sigma_{\varepsilon}^{+})) F_{\varepsilon}(S_{\varepsilon} v_{\varepsilon}(s), \Sigma_{\varepsilon}(v_{\varepsilon}(s))) ds, \quad \varepsilon \in [0, 1].$$
(6.5)

in the class of Lipschitz maps $\Sigma_{\varepsilon}: \mathbb{R}^n \to \operatorname{Ker}(Q_{\varepsilon})$ which are globally bounded with bound D and globally Lipschitz with Lipschitz constant L.

Note that, from (4.13),

$$\|\Phi(\Sigma_{\varepsilon})(\eta)\|_{U_{\varepsilon}^{p}} = \int_{-\infty}^{\tau} \rho C(\tau - s)^{-\gamma} e^{-b(\tau - s)} ds.$$
 (6.6)

and for suitably chosen ρ we have that $\|\Phi(\Sigma_{\varepsilon})\| \leq D$.

Next, suppose that Σ_{ε} and $\tilde{\Sigma}_{\varepsilon}$ are functions satisfying (6.4), η , $\tilde{\eta} \in \mathbb{R}^n$ and denote $v_{\varepsilon}(t) = \psi(t, \tau, \eta, \Sigma_{\varepsilon})$, $\tilde{v}_{\varepsilon}(t) = \psi(t, \tau, \tilde{\eta}, \tilde{\Sigma}_{\varepsilon})$. Then,

$$v_{\varepsilon}(t) - \tilde{v}_{\varepsilon}(t) = e^{-B_{\varepsilon}(t-\tau)}(\eta - \tilde{\eta}) + \int_{\tau}^{t} e^{-B_{\varepsilon}(t-s)}Q_{\varepsilon}[F_{\varepsilon}(S_{\varepsilon}v_{\varepsilon}, \Sigma_{\varepsilon}(v_{\varepsilon})) - F_{\varepsilon}(S_{\varepsilon}\tilde{v}_{\varepsilon}, \tilde{\Sigma}_{\varepsilon}(\tilde{v}_{\varepsilon}))]ds.$$

And

$$\begin{split} \|v_{\varepsilon}(t) - \tilde{v_{\varepsilon}}(t)\|_{\mathbb{R}^{n}} &\leqslant Ce^{b(t-\tau)} \|\eta - \tilde{\eta}\|_{\mathbb{R}^{n}} \\ &+ C\int_{t}^{\tau} e^{b(t-s)} \|Q_{\varepsilon}F_{\varepsilon}(S_{\varepsilon}v_{\varepsilon}, \Sigma_{\varepsilon}(v_{\varepsilon})) - Q_{\varepsilon}F_{\varepsilon}(S_{\varepsilon}\tilde{v}_{\varepsilon}, \tilde{\Sigma}_{\varepsilon}(\tilde{v}_{\varepsilon}))\|_{\mathbb{R}^{n}} ds \\ &\leqslant Ce^{b(t-\tau)} \|\eta - \tilde{\eta}\|_{\mathbb{R}^{n}} \\ &+ \rho C\int_{t}^{\tau} e^{-b(t-s)} \left(\|\Sigma_{\varepsilon}(v_{\varepsilon}) - \tilde{\Sigma}_{\varepsilon}(\tilde{v}_{\varepsilon})\|_{U_{\varepsilon}^{p}} + \|v_{\varepsilon} - \tilde{v}_{\varepsilon}\|_{\mathbb{R}^{n}} \right) ds \\ &\leqslant Ce^{b(t-\tau)} \|\eta - \tilde{\eta}\|_{\mathbb{R}^{n}} \\ &+ \rho C\int_{t}^{\tau} e^{b(t-s)} \left(\|\Sigma_{\varepsilon}(\tilde{v}_{\varepsilon}) - \tilde{\Sigma}_{\varepsilon}(\tilde{v}_{\varepsilon})\|_{U_{\varepsilon}^{p}} + (1+L)\|v_{\varepsilon} - \tilde{v}_{\varepsilon}\|_{\mathbb{R}^{n}} \right) ds \\ &\leqslant Ce^{b(t-\tau)} \|\eta - \tilde{\eta}\|_{\mathbb{R}^{n}} \\ &+ \rho C\int_{t}^{\tau} e^{b(t-s)} \left((1+L) \|v_{\varepsilon} - \tilde{v}_{\varepsilon}\|_{\mathbb{R}^{n}} + \|\Sigma_{\varepsilon} - \tilde{\Sigma}_{\varepsilon}\|_{U_{\varepsilon}^{p}} \right) ds \\ &\leqslant Ce^{b(t-\tau)} \|\eta - \tilde{\eta}\|_{\mathbb{R}^{n}} \\ &+ \rho C(1+L) \int_{t}^{\tau} e^{b(t-s)} \|v_{\varepsilon} - \tilde{v}_{\varepsilon}\|_{\mathbb{R}^{n}} ds + \rho C(\|\Sigma_{\varepsilon} - \tilde{\Sigma}_{\varepsilon}\|_{U_{\varepsilon}^{p}} \int_{t}^{\tau} e^{b(t-s)} ds. \end{split}$$

Let $\phi(t) = e^{-b(t-\tau)} \|v_{\varepsilon}(t) - \tilde{v_{\varepsilon}}(t)\|_{\mathbb{R}^n}$. Then,

$$\phi(t) \leqslant C \|\eta - \tilde{\eta}\|_{\mathbb{R}^n} + \rho C \int_t^{\tau} e^{b(\tau - s)} ds |\|\Sigma_{\varepsilon} - \tilde{\Sigma}_{\varepsilon}|\|_{U_{\varepsilon}^p} + C \rho (1 + L) \int_t^{\tau} \phi(s) ds.$$

By Gronwall's inequality

$$||v_{\varepsilon}(t) - \tilde{v_{\varepsilon}}(t)||_{\mathbb{R}^{n}} \leq |C||\eta - \tilde{\eta}||_{\mathbb{R}^{n}} e^{b(t-\tau)} + \rho C \int_{t}^{\tau} e^{b(t-s)} ds ||\Sigma_{\varepsilon} - \tilde{\Sigma}_{\varepsilon}||_{U_{\varepsilon}^{p}} e^{-\rho C(1+L)(t-\tau)}$$
$$\leq |C||\eta - \tilde{\eta}||_{\mathbb{R}^{n}} + \rho C b^{-1} ||\Sigma_{\varepsilon} - \tilde{\Sigma}_{\varepsilon}||_{U_{\varepsilon}^{p}} e^{-\rho C(1+L)(t-\tau)}$$

Thus,

$$\begin{split} \|\Phi(\Sigma_{\varepsilon})(\eta) - \Phi(\tilde{\Sigma}_{\varepsilon})(\tilde{\eta})\|_{U_{\varepsilon}^{p}} \\ &\leqslant C \int_{-\infty}^{\tau} (\tau - s)^{-\gamma} e^{-b(\tau - s)} \|F_{\varepsilon}(S_{\varepsilon}v_{\varepsilon}, \Sigma_{\varepsilon}(v_{\varepsilon})) - F_{\varepsilon}(S_{\varepsilon}\tilde{v}_{\varepsilon}, \tilde{\Sigma}_{\varepsilon}(\tilde{v}_{\varepsilon}))\|_{L^{2}(\Omega_{\varepsilon})} ds \\ &\leqslant \rho C \int_{-\infty}^{\tau} (\tau - s)^{-\gamma} e^{-b(\tau - s)} \left(\|\Sigma_{\varepsilon}(v_{\varepsilon}) - \tilde{\Sigma}_{\varepsilon}(\tilde{v}_{\varepsilon})\|_{U_{\varepsilon}^{p}} + \|v_{\varepsilon} - \tilde{v}_{\varepsilon}\|_{\mathbb{R}^{n}} \right) ds \\ &\leqslant \rho C \int_{-\infty}^{\tau} (\tau - s)^{-\gamma} e^{-b(\tau - s)} \left[(1 + L) \|v_{\varepsilon} - \tilde{v}_{\varepsilon}\|_{\mathbb{R}^{n}} + |\|\Sigma_{\varepsilon} - \tilde{\Sigma}_{\varepsilon}\|\| \right] ds. \end{split}$$

Using the estimates for $||v_{\varepsilon} - \tilde{v}_{\varepsilon}||_{\mathbb{R}^n}$ we obtain

$$\|\Phi(\Sigma_{\varepsilon})(\eta) - \Phi(\tilde{\Sigma}_{\varepsilon})(\tilde{\eta})\| \leqslant \rho C \Gamma(1-\gamma) \left[b^{-1+\gamma} + \frac{\rho C(1+L)}{b(b-\rho C(1+L))^{1-\gamma}} \right] \|\Sigma_{\varepsilon} - \tilde{\Sigma}_{\varepsilon}\|$$

$$+ \frac{\rho C^{2}(1+L)\Gamma(1-\gamma)}{(b-\rho C(1+L))^{-1+\gamma}} \|\eta - \tilde{\eta}\|_{\mathbb{R}^{n}}.$$

Let

$$I_{\Sigma}(\rho) = \rho C \Gamma(1 - \gamma) \left[b^{-1+\gamma} + \frac{\rho C(1+L)}{b(b-\rho C(1+L))^{1-\gamma}} \right]$$

and

$$I_{\eta}(\rho) = \frac{\rho C^{2}(1+L)\Gamma(1-\gamma)}{(b-\rho C(1+L))^{1-\gamma}}.$$

It is easy to see that, given $\theta < 1$, there exists a ρ_0 such that, for $\rho \leqslant \rho_0$, $I_{\Sigma}(\rho) \leqslant \theta$ and $I_{\eta}(\rho) \leqslant L$ and

$$\|\Phi(\Sigma_{\varepsilon})(\eta) - \Phi(\tilde{\Sigma}_{\varepsilon})(\tilde{\eta})\|_{U_{\varepsilon}^{p}} \leqslant L\|\eta - \eta'\|_{\mathbb{R}^{n}} + \theta \|\Sigma_{\varepsilon} - \tilde{\Sigma}_{\varepsilon}\|.$$

$$(6.7)$$

The inequalities (6.6) and (6.7) imply that G is a contraction map from the class of functions that satisfy (6.4) into itself. Therefore, it has a unique fixed point $\Sigma_{\varepsilon}^* = \Phi(\Sigma_{\varepsilon}^*)$ in this class. The invariance follows in the usual manner.

The fact that the graph is the whole unstable manifold follows (taking the limit as t_0 tends to $-\infty$) from the following: If $w(t) = (v(t), z(t)), t \in \mathbb{R}$, is a global solution of (6.1) which is bounded as $t \to -\infty$, there are constants $\tilde{M} \geq 1$ and $\nu > 0$ such that

$$||z(t) - \Sigma_{\varepsilon}^*(v(t))||_{U_{\varepsilon}^p} \leqslant \tilde{M}(t - t_0)^{-\gamma} e^{-\nu(t - t_0)} ||z(t_0) - \Sigma_{\varepsilon}^*(v(t_0))||_{U_{\varepsilon}^p}, \quad t_0 < t.$$
(6.8)

The proof of (6.8) can be carried out following the steps in the proof of (A.8) in [6], using the singular Gronwall's inequality instead of the usual one, and noting that ε can be considered fixed for this purpose.

It remains to prove the continuity of the unstable manifolds. This is accomplished in the following manner. If $0 \le \varepsilon \le \varepsilon_0$ is such that the unstable manifold is given by the graph of Σ_{ε}^* , $0 \le \varepsilon \le \varepsilon_0$, we want to show that

$$\sup_{\eta \in \mathbb{R}^n} \|\Sigma_{\varepsilon}^*(\eta) - E_{\varepsilon} \Sigma_0^*(\eta)\|_{U_{\varepsilon}^p} = \|\Sigma_{\varepsilon}^* - E_{\varepsilon} \Sigma_0^*\|.$$

It follows from Proposition 4.3 that

$$\begin{split} &\|\Sigma_{\varepsilon}^{*}(\eta_{\varepsilon}) - E_{\varepsilon}\Sigma_{0}^{*}(\eta)\|_{U_{\varepsilon}^{p}} \\ &\leq \int_{-\infty}^{\tau} \|e^{-\tilde{A}_{\varepsilon}(\tau-s)}(I - Q(\sigma_{\varepsilon}^{+}))F_{\varepsilon}(S_{\varepsilon}v_{\varepsilon}, \Sigma_{\varepsilon}^{*}(v_{\varepsilon})) - E_{\varepsilon}e^{-\tilde{A}_{0}(\tau-s)}(I - Q(\sigma_{0}^{+}))F_{0}(S_{0}v_{0}, \Sigma_{0}^{*}(v_{0}))\|_{U_{\varepsilon}^{p}} ds \\ &\leq \int_{-\infty}^{\tau} \|e^{-\tilde{A}_{\varepsilon}(\tau-s)}(I - Q(\sigma_{\varepsilon}^{+}))[F_{\varepsilon}(S_{\varepsilon}v_{\varepsilon}, \Sigma_{\varepsilon}^{*}(v_{\varepsilon})) - E_{\varepsilon}F_{0}(S_{0}v_{0}, \Sigma_{0}^{*}(v_{0}))]\|_{U_{\varepsilon}^{p}} ds \\ &+ \int_{-\infty}^{\tau} \|[e^{-\tilde{A}_{\varepsilon}(\tau-s)}(I - Q(\sigma_{\varepsilon}^{+}) - E_{\varepsilon}e^{-\tilde{A}_{0}(\tau-s)}(I - Q(\sigma_{0}^{+})M_{\varepsilon}]E_{\varepsilon}F_{0}(S_{0}v_{0}, \Sigma_{0}^{*}(v_{0}))\|_{U_{\varepsilon}^{p}} ds \\ &\leq \int_{-\infty}^{\tau} \|e^{-\tilde{A}_{\varepsilon}(\tau-s)}(I - Q(\sigma_{\varepsilon}^{+}))[F_{\varepsilon}(S_{\varepsilon}v_{\varepsilon}, \Sigma_{\varepsilon}^{*}(v_{\varepsilon})) - F_{\varepsilon}(E_{\varepsilon}(S_{\varepsilon}v_{0}, \Sigma_{0}^{*}(v_{0})))]\|_{U_{\varepsilon}^{p}} ds \\ &+ \int_{-\infty}^{\tau} \|[e^{-\tilde{A}_{\varepsilon}(\tau-s)}(I - Q(\sigma_{\varepsilon}^{+}) - E_{\varepsilon}e^{\tilde{A}_{0}(\tau-s)}(I - Q(\sigma_{0}^{+})M_{\varepsilon}]E_{\varepsilon}F_{0}(S_{0}v_{0}, \Sigma_{0}^{*}(v_{0})))\|_{U_{\varepsilon}^{p}} ds \\ &\leq C\int_{-\infty}^{\tau} e^{b(\tau-s)}(T - S)^{-\gamma}\|F_{\varepsilon}(S_{\varepsilon}v_{\varepsilon}, \Sigma_{\varepsilon}^{*}(v_{\varepsilon})) - F_{\varepsilon}(E_{\varepsilon}(S_{0}v_{0}, \Sigma_{0}^{*}(v_{0})))\|_{U_{\varepsilon}^{p}} ds \\ &+ C\rho(\varepsilon)\int_{-\infty}^{\tau} e^{b(\tau-s)}\|F_{\varepsilon}(S_{\varepsilon}v_{\varepsilon}, \Sigma_{\varepsilon}^{*}(v_{\varepsilon}))\|_{C(\bar{\Omega}_{\varepsilon})} ds \\ &\leq \rho Cb^{-1}\rho(\varepsilon) + \rho Cb^{\gamma-1}\Gamma(1 - \gamma)\|\Sigma_{\varepsilon}^{*} - E_{\varepsilon}\Sigma_{0}^{*}\|\| \\ &+ \rho C(1 + L)\int_{-\infty}^{\tau} e^{-b(\tau-s)}(\tau - s)^{-\gamma}\|v_{\varepsilon} - v_{0}\|_{\mathbb{R}^{n}} ds. \end{split}$$

Thus, it is enough to estimate $||v_{\varepsilon} - v_0||_{\mathbb{R}^n}$. Note that

$$\begin{split} \|v_{\varepsilon} - v_{0}\|_{\mathbb{R}^{n}} & \leqslant \int_{t}^{\tau} \|e^{-B_{\varepsilon}(t-s)} - e^{-B_{0}(t-s)}\| \|F_{\varepsilon}(S_{\varepsilon}v_{\varepsilon}, \Sigma_{\varepsilon}^{*}(v_{\varepsilon}))\|_{\mathbb{R}^{n}} ds \\ & + \int_{t}^{\tau} \|e^{-B_{0}(t-s)}\| \|F_{\varepsilon}(S_{\varepsilon}v_{\varepsilon}, \Sigma_{\varepsilon}^{*}(v_{\varepsilon})) - F_{0}(S_{0}v_{0}, \Sigma_{0}^{*}(v_{0}))\|_{\mathbb{R}^{n}} ds \\ & \leqslant \rho M b^{-1}[o(1) + |\|\Sigma_{\varepsilon}^{*} - \Sigma_{0}^{*}|\|] + \rho C(1 + L) \int_{t}^{\tau} e^{b(t-s)} \|v_{\varepsilon} - v_{0}\|_{\mathbb{R}^{n}} ds \end{split}$$

Therefore

$$||v_{\varepsilon} - v_0||_{\mathbb{R}^n} \le \rho C b^{-1} [o(1) + |||\Sigma_{\varepsilon}^* - \Sigma_0^*|||] e^{-\rho C(1+L)(\tau-t)}$$

which shows that

$$\sup_{\eta \in \mathbb{R}^n} \|\Sigma_{\varepsilon}^*(\eta) - \Sigma_0^*(\eta)\|_{U_{\varepsilon}^p} \stackrel{\varepsilon \to 0}{\to} 0.$$

This proves the result.

7. Continuity of Attractors in Other Norms

In this section we study the continuity of attractors in other norms and very specially in the norm of the space $U_{\varepsilon}^{1,2}$, see (2.8). This continuity is obtained as a consequence of the regularization properties of the nonlinear semigroups. As a matter of fact, in many instances the attractors $\mathcal{A}_{\varepsilon}$, \mathcal{A}_{0} live in better spaces X_{ε} and X_{0} respectively for which the linear map $E_{\varepsilon}: X_0 \to X_{\varepsilon}$ is well defined as well. We would like to give conditions that, once the continuity of the attractors in U_{ε}^p is obtained, will guarantee the continuity results for the attractors in these better spaces. In fact, the following result holds

Proposition 7.1. If there exists a $\tau > 0$ fixed such that for each sequence of $\varepsilon_n \to 0$, $\phi_{\varepsilon_n} \in \mathcal{A}_{\varepsilon_n}$ and $\phi_0 \in \mathcal{A}_0$ with $\|\phi_{\varepsilon_n} - E_{\varepsilon_n}\phi_0\|_{U^p_{\varepsilon_n}} \to 0$ implies that

$$||T_{\varepsilon_n}(\tau,\phi_{\varepsilon_n}) - E_{\varepsilon_n}T_0(\tau,\phi_0)||_{X_{\varepsilon_n}} \to 0$$
(7.1)

then, the upper semicontinuity of the attractors in U^p_{ε} implies the upper semicontinuity in X_{ε} and the lower semicontinuity of the attractors in U^p_{ε} implies the lower semicontinuity of the attractors in X_{ε} .

Proof: Assume we have a family of $\varphi_{\varepsilon} \in \mathcal{A}_{\varepsilon}$. From the invariance of the attractors under the semigroup T_{ε} , we have that there exist $\phi_{\varepsilon} \in \mathcal{A}_{\varepsilon}$ with $T_{\varepsilon}(\tau, \phi_{\varepsilon}) = \varphi_{\varepsilon}$.

If the attractors are E_{ε} -upper semicontinuous in U_{ε}^{p} , we have that for each sequence $\varepsilon_{n} \to 0$, there will exist a subsequence, that we still denote by ε_{n} and an element $\phi_{0} \in \mathcal{A}_{0}$ such that $\|\phi_{\varepsilon_{n}} - E_{\varepsilon_{n}}\phi_{0}\|_{U_{\varepsilon_{n}}^{p}} \to 0$ as $\varepsilon_{n} \to 0$. With (7.1) we get that if we define $\varphi_{0} = T_{0}(\tau, \phi_{0})$, we have $\|\varphi_{\varepsilon_{n}} - E_{\varepsilon_{n}}\varphi_{0}\|_{X_{\varepsilon_{n}}} \to 0$, which shows the E_{ε} -upper semicontinuity in X_{ε} .

Assume now that the attractors are E_{ε} -lower semicontinuous in U_{ε}^{p} . If $\varphi_{0} \in \mathcal{A}_{0}$ and if we define $\phi_{0} \in \mathcal{A}_{0}$ with $T_{0}(\tau, \phi_{0}) = \varphi_{0}$, then there will exist a sequence of $\phi_{\varepsilon} \in \mathcal{A}_{\varepsilon}$ with $\|\phi_{\varepsilon} - E_{\varepsilon}\phi_{0}\|_{U_{\varepsilon}^{p}} \to 0$ as $\varepsilon \to 0$. Using (7.1) again, we get that $\|\varphi_{\varepsilon} - E_{\varepsilon}\varphi_{0}\|_{X_{\varepsilon}} \to 0$ which shows the E_{ε} -lower semicontinuty in X_{ε} .

With this result we can provide now a proof of Theorem 2.7.

Proof of Theorem 2.7: We will apply Proposition 7.1, proving first that

$$||T_{\varepsilon_n}(\tau,\phi_{\varepsilon_n}) - E_{\varepsilon_n}T_0(\tau,\phi_0)||_{U_{\varepsilon_n}^{1,2}} \to 0$$

for some $\tau > 0$ fixed, sequences $\varepsilon_n \to 0$, $\phi_{\varepsilon_n} \in \mathcal{A}_{\varepsilon_n}$ and $\phi_0 \in \mathcal{A}_0$ with $\|\phi_{\varepsilon_n} - E_{\varepsilon_n}\phi_0\|_{U^p_{\varepsilon_n}} \to 0$. Observe first that

$$||T_{\varepsilon_n}(\tau,\phi_{\varepsilon_n}) - E_{\varepsilon_n}T_0(\tau,\phi_0)||_{U_{\varepsilon_n}^{1,2}} \leq ||T_{\varepsilon_n}(\tau,\phi_{\varepsilon_n}) - E_{\varepsilon_n}T_0(\tau,M_{\varepsilon}\phi_{\varepsilon_n})||_{U_{\varepsilon_n}^{1,2}} + ||E_{\varepsilon}T_0(\tau,M_{\varepsilon}\phi_{\varepsilon_n}) - E_{\varepsilon_n}T_0(\tau,\phi_0)||_{U_{\varepsilon_n}^{1,2}}$$
(7.2)

and for a fixed $\tau > 0$,

$$||E_{\varepsilon}T_0(\tau, M_{\varepsilon}\phi_{\varepsilon_n}) - E_{\varepsilon_n}T_0(\tau, \phi_0)||_{U_{\varepsilon_n}^{1,2}} \leq ||T_0(\tau, M_{\varepsilon}\phi_{\varepsilon_n}) - T_0(\tau, \phi_0)||_{U_0^{1,2}} \to 0$$

since $T_0(\tau,\cdot):U_0^p\to U_0^{1,2}$ is continuous, see [4] .

To estimate the first term of the second line of (7.2) we use the Variation of Constants Formula (5.1) for $\varepsilon \in [0,1]$ and with simple computations we obtain

$$||T_{\varepsilon}(t,\phi_{\varepsilon}) - E_{\varepsilon}T_{0}(t,M_{\varepsilon}\phi_{\varepsilon})||_{U_{\varepsilon}^{1,2}} \leq ||e^{-A_{\varepsilon}t}\phi_{\varepsilon} - E_{\varepsilon}e^{-A_{0}t}M_{\varepsilon}\phi_{\varepsilon}||_{U_{\varepsilon}^{1,2}} + \int_{0}^{t} ||\left(e^{-A_{\varepsilon}(t-s)} - E_{\varepsilon}e^{-A_{0}(t-s)}M_{\varepsilon}\right)f_{\varepsilon}(T_{\varepsilon}(s,\phi_{\varepsilon}))||_{U_{\varepsilon}^{1,2}}ds + \int_{0}^{t} ||E_{\varepsilon}e^{-A_{0}(t-s)}M_{\varepsilon}(f_{\varepsilon}(T_{\varepsilon}(s,\phi_{\varepsilon})) - f_{\varepsilon}(E_{\varepsilon}T_{0}(s,M_{\varepsilon}\phi_{\varepsilon})))||_{U_{\varepsilon}^{1,2}}ds, \quad \varepsilon \in [0,\varepsilon_{0}].$$

$$(7.3)$$

But note that $\mathcal{A}_{\varepsilon} \subset C(\bar{\Omega}_{\varepsilon})$ for $0 < \varepsilon \leqslant \varepsilon_0$, $\mathcal{A}_0 \subset C(\bar{\Omega}) \oplus C([0,1])$ and that we have uniform bounds in these spaces.

If we are able to obtain the following two estimates:

$$\|e^{-A_{\varepsilon}t} - E_{\varepsilon}e^{-A_{0}t}M_{\varepsilon}\|_{\mathcal{L}(C(\bar{\Omega}) \oplus C(\bar{R}_{\varepsilon}), U_{\varepsilon}^{1,2})} \leqslant Ct^{-\gamma}\nu(\varepsilon), \ t > 0.$$

$$(7.4)$$

for some $0 \leqslant \gamma < 1$ and with $\nu(\varepsilon) \to 0$ as $\varepsilon \to 0$, and

$$||e^{-A_0t}||_{\mathcal{L}(U^p_t, U^{1,2}_0)} \le Ct^{-\beta}, \quad t > 0.$$
 (7.5)

for some $0 \le \beta < 1$, then using (7.4) and (7.5) in (7.3) and using the convergence of the nonlinear semigroup in U_0^p we obtain that $||T_{\varepsilon}(t,\phi_{\varepsilon}) - E_{\varepsilon}T_0(t,M_{\varepsilon}\phi_{\varepsilon})||_{U_{\varepsilon}^{1,2}} \to 0$ as $\varepsilon \to 0$.

The proof of (7.5) is in [4] Remark 3.2.

Hence we just need to show (7.4). To obtain this estimate we need some extra resolvent estimates, similar to the ones obtained in Section 3.1. To that end we introduce the continuous extension operator

$$E_{\varepsilon}^{\mathcal{C}}: \quad \mathcal{C}(\overline{\Omega}) \oplus \mathcal{C}(0,1) \quad \to \quad \mathcal{C}(\overline{\Omega_{\varepsilon}})$$

$$(w_{\varepsilon}, v_{\varepsilon}) \quad \to \quad E_{\varepsilon}^{\mathcal{C}}(w_{\varepsilon}, v_{\varepsilon}) = \begin{cases} w_{\varepsilon}, & x \in \Omega \\ \tilde{v}_{\varepsilon}, & x \in \mathbf{R}_{\varepsilon}, \end{cases}$$

$$(7.6)$$

where

 $\tilde{v}_{\varepsilon}(x) = v_{\varepsilon}(s) + h_{\varepsilon}(s) \left(w_{\varepsilon}(0, y) - v_{\varepsilon}(0)\right) + h_{\varepsilon}(1 - s) \left(w_{\varepsilon}(1, y) - v_{\varepsilon}(1)\right), \ x = (s, y) \in \mathbb{R}_{\varepsilon},$ (7.7) and the function $h_{\delta}(s) = h(\frac{s}{\delta})$, where $h : \mathbb{R}^+ \to [0, 1]$ is a \mathcal{C}^{∞} function such that

$$h(s) = \begin{cases} 1, & \text{for } s \in [0, 1/4], \\ 0, & \text{for } s \geqslant 3/4 \end{cases}$$

and $|h'(s)| \leq C$.

Observe that with this definition $E_{\varepsilon}^{\mathcal{C}}(w_{\varepsilon}, v_{\varepsilon})$ is always a continuous function in $\bar{\Omega}_{\varepsilon}$ if $(w_{\varepsilon}, v_{\varepsilon}) \in \mathcal{C}(\overline{\Omega}) \oplus \mathcal{C}(0, 1)$. Moreover, if $(w_{\varepsilon}, v_{\varepsilon}) \in U_0^{1,2}$ then, $E_{\varepsilon}^{\mathcal{C}}(w_{\varepsilon}, v_{\varepsilon}) \in H^1(\Omega_{\varepsilon})$.

We also need the following lemmas whose proofs will be provided later.

Lemma 7.2. Let $\lambda \in \rho(A_{\varepsilon}) \cap \rho(A_0)$, then the following holds

$$(\lambda + A_{\varepsilon})^{-1} - E_{\varepsilon}(\lambda + A_{0})^{-1}M = (I - \lambda(A_{\varepsilon} + \lambda)^{-1})(A_{\varepsilon}^{-1} - E_{\varepsilon}^{\mathcal{C}}A_{0}^{-1}M_{\varepsilon})(I - \lambda E_{\varepsilon}(A_{0} + \lambda)^{-1}M_{\varepsilon}) + (I - \lambda(A_{\varepsilon} + \lambda)^{-1})(E_{\varepsilon}^{\mathcal{C}} - E_{\varepsilon})(A_{0} + \lambda)^{-1}M_{\varepsilon}$$

Lemma 7.3. There is a constant C > 0 such that for each $\lambda \in \Sigma_{\theta}$ we have

$$\|(E_{\varepsilon}^{\mathcal{C}} - E_{\varepsilon})(A_0 + \lambda)^{-1} M_{\varepsilon}\|_{\mathcal{L}(\mathcal{C}(\overline{\Omega}) \oplus \mathcal{C}(\bar{R}_{\varepsilon}), H^1(\Omega) \oplus H^1(R_{\varepsilon}))} \leqslant C \frac{\varepsilon^{\frac{N}{2}}}{1 + |\lambda|^{\frac{1}{2}}},$$

$$\|(I - \lambda(A_{\varepsilon} + \lambda)^{-1})(E_{\varepsilon}^{\mathcal{C}} - E_{\varepsilon})(A_0 + \lambda)^{-1} M_{\varepsilon}\|_{\mathcal{L}(\mathcal{C}(\overline{\Omega}) \oplus \mathcal{C}(\bar{R}_{\varepsilon}), H^1(\Omega) \oplus H^1(R_{\varepsilon}))} \leqslant C\varepsilon^{N/2}.$$

$$(7.8)$$

Lemma 7.4. There is a constant C > 0, independent of ε , such that

$$(i) \|E_{\varepsilon}(I - \lambda(A_0 + \lambda)^{-1})M_{\varepsilon}f_{\varepsilon}\|_{C(\bar{\Omega}) \oplus C(\bar{R}_{\varepsilon})} \leqslant C \|f_{\varepsilon}\|_{C(\bar{\Omega}) \oplus C(\bar{R}_{\varepsilon})},$$

$$(ii) \|(I - \lambda(A_{\varepsilon} + \lambda)^{-1})g_{\varepsilon}\|_{H^{1}(\Omega_{\varepsilon})} \leq C \|g_{\varepsilon}\|_{H^{1}(\Omega_{\varepsilon})}.$$

Lemma 7.5. There exists a constant C > 0 such that for all $\lambda \in \Sigma_{\theta}$ and all $f_{\varepsilon} \in \mathcal{C}(\overline{\Omega}) \oplus$ $\mathcal{C}(R_{\varepsilon}),$

$$\left\| ((A_{\varepsilon} + \lambda)^{-1} - E_{\varepsilon}(A_0 + \lambda)^{-1} M_{\varepsilon}) f_{\varepsilon} \right\|_{H^{1}(\Omega) \oplus H^{1}(R_{\varepsilon})} \leqslant C \varepsilon^{N/2} \| f_{\varepsilon} \|_{\mathcal{C}(\overline{\Omega}) \oplus \mathcal{C}(R_{\varepsilon})}.$$
 (7.9)

Clearly, from Lemma 7.5 and the expression of the differences of the semigroups in terms of the integral of the difference of the resolvents as in (4.9), we have that there is a constant C > 0 such that

$$\left\| e^{-A_{\varepsilon}t} - E_{\varepsilon}e^{-A_{0}t}M_{\varepsilon} \right\|_{\mathcal{L}(C(\bar{\Omega}) \oplus C(\bar{R}_{\varepsilon}), H^{1}(\Omega) \oplus H^{1}(R_{\varepsilon}))} \leqslant C\varepsilon^{N/2}t^{-1}.$$

$$(7.10)$$

On the other hand

$$\|e^{-A_{\varepsilon}t} - E_{\varepsilon}e^{-A_{0}t}M_{\varepsilon}\|_{\mathcal{L}(L^{\infty}(\Omega_{\varepsilon}),H^{1}(\Omega)\oplus H^{1}(R_{\varepsilon}))} \le$$

$$\left\| e^{-A_{\varepsilon}t} \right\|_{\mathcal{L}(L^{\infty}(\Omega_{\varepsilon}), H^{1}(\Omega) \oplus H^{1}(R_{\varepsilon}))} + \left\| E_{\varepsilon}e^{-A_{0}t}M_{\varepsilon} \right\|_{\mathcal{L}(L^{\infty}(\Omega_{\varepsilon}), H^{1}(\Omega) \oplus H^{1}(R_{\varepsilon}))} \le \tag{7.11}$$

$$\|e^{-A_{\varepsilon}t}\|_{\mathcal{L}(L^{2}(\Omega_{\varepsilon}),H^{1}(\Omega_{\varepsilon}))} + \|e^{-A_{0}t}\|_{\mathcal{L}(U_{\varepsilon}^{p},H^{1}(\Omega)\oplus H^{1}(0,1))} \le Ct^{-\beta}$$

for some β with $1/2 < \beta < 1$, see [4], Remark 3.2. Interpolating (7.10) and (7.11), we have that that, for any $\eta < 1$,

$$\left\| e^{-A_{\varepsilon}t} - E_{\varepsilon}e^{-A_{0}t}M_{\varepsilon} \right\|_{\mathcal{L}(L^{\infty}(\Omega_{\varepsilon}), H^{1}(\Omega) \oplus H^{1}(R_{\varepsilon}))} \leqslant C\varepsilon^{\eta N/2}t^{-(\eta + (1-\eta)\beta)}. \tag{7.12}$$

Choosing $\frac{N-1}{N} < \eta < 1$ so that $\eta N/2 > (N-1)/2$, the result follows with $\gamma = \eta + (1-\eta)\beta < 1$. This shows estimate (7.4) and the theorem is proved.

Remark 7.6. We may also obtain the convergence of the attractors in some other norms. As a matter of fact if K is a compact subset of $\bar{\Omega} \setminus \{P_0, P_1\}$ we can easily obtain uniform bounds of all the attractors for instance in $C^{1,\eta}(K)$. This estimates may be obtained with an appropriate cut-off function and using standard regularity properties of the nonlinear semigroups (we are far away from the channel R_{ε}). Hence, since we have obtained already the continuity (lower or upper) of the attractor in $L^p(K)$, with the compact embedding of $C^{1,\eta}(K)$ in $C^{1,\eta^-}(K)$ we also get the continuity (lower or upper) in $C^{1,\eta^-}(K)$.

We provide now the proofs of the different lemmas we have stated above.

Proof of Lemma 7.2: This lemma is obtained in a similar way as Lemma 3.5.

Proof of Lemma 7.3: Let $f_{\varepsilon} \in \mathcal{C}(\overline{\Omega}) \oplus \mathcal{C}(\overline{R}_{\varepsilon})$ and define $K_{\varepsilon} := (E_{\varepsilon}^{\mathcal{C}} - E_{\varepsilon})(A_0 + \lambda)^{-1}Mf_{\varepsilon} = \tilde{z}_{\varepsilon} - z_{\varepsilon}$,, where $\tilde{z}_{\varepsilon} = E_{\mathcal{C}}(A_0 + \lambda)^{-1}Mf_{\varepsilon}$ e $z_{\varepsilon} = E_{\varepsilon}(A_0 + \lambda)^{-1}Mf_{\varepsilon}$. Observe that $(A_0 + \lambda)^{-1}Mf_{\varepsilon} = (w_{\varepsilon}, v_{\varepsilon})$ where

$$\begin{cases}
-\Delta w_{\varepsilon} + \lambda w_{\varepsilon} = f_{\varepsilon}, & x \in \Omega \\
\frac{\partial w_{\varepsilon}}{\partial n} = 0, & x \in \partial \Omega \\
-\frac{1}{g} (gv_{\varepsilon s})_{s} + \lambda v_{\varepsilon} = M f_{\varepsilon}, & s \in (0, 1) \\
v_{\varepsilon}(0) = w_{\varepsilon}(0), v_{\varepsilon}(1) = w_{\varepsilon}(1),
\end{cases} (7.13)$$

 $\tilde{v}_{\varepsilon}(s,y) = v_{\varepsilon}(s) + h_{\varepsilon}(s) \left(w_{\varepsilon}(0,y) - v_{\varepsilon}(0)\right) + h_{\varepsilon}(1-s) \left(w_{\varepsilon}(1,y) - v_{\varepsilon}(1)\right), \ \forall (s,y) \in R_{\varepsilon} \text{ and }$ $z_{\varepsilon}(s,y) = v_{\varepsilon}(s), \forall (s,y) \in R_{\varepsilon}.$

Also note that since $K_{\varepsilon} \equiv 0$ in Ω , we have $\|K_{\varepsilon}\|_{H^{1}(\Omega) \oplus H^{1}(R_{\varepsilon})} = \|K_{\varepsilon}\|_{H^{1}(R_{\varepsilon})}^{2}$. Moreover,

$$||K_{\varepsilon}||_{L^{2}(R_{\varepsilon})}^{2} \leq 2 \int_{0}^{\varepsilon} \int_{\Gamma_{\varepsilon}^{s}} |h_{\varepsilon}(s)|^{2} |w_{\varepsilon}(0, y) - v_{\varepsilon}(0)|^{2} ds dy$$

$$+ 2 \int_{1-\varepsilon}^{1} \int_{\Gamma_{\varepsilon}^{s}} |h_{\varepsilon}(1-s)|^{2} |w_{\varepsilon}(1, y) - v_{\varepsilon}(1)|^{2} ds dy$$

$$\leq C_{2} \varepsilon^{N} ||w_{\varepsilon}||_{\mathcal{C}(\overline{\Omega})}^{2}.$$

Now note that $h'_{\varepsilon}(s) = \varepsilon^{-1}h'(x/\varepsilon)$, $h'_{\varepsilon}(1-s) = -\varepsilon^{-1}h'((1-s)/\varepsilon)$. Hence, with similar estimates as above,

$$\begin{split} \|\nabla K_{\varepsilon}\|_{L^{2}(R_{\varepsilon})}^{2} &\leqslant 2 \int_{0}^{\varepsilon} |h_{\varepsilon}'(s)|^{2} \int_{\Gamma_{\varepsilon}^{s}} |w_{\varepsilon}(0,y) - v_{\varepsilon}(0)|^{2} \, ds dy \\ &+ 2 \int_{1-\varepsilon}^{1} |h_{\varepsilon}'(1-s)|^{2} \int_{\Gamma_{\varepsilon}^{s}} |w_{\varepsilon}(1,y) - v_{\varepsilon}(1)|^{2} \, ds dy \\ &+ 2 \int_{0}^{\varepsilon} |h_{\varepsilon}(s)|^{2} \int_{\Gamma_{\varepsilon}^{s}} |\nabla_{y} w_{\varepsilon}(0,y)|^{2} \, ds dy + 2 \int_{1-\varepsilon}^{1} |h_{\varepsilon}(1-s)|^{2} \int_{\Gamma_{\varepsilon}^{s}} |\nabla_{y} w_{\varepsilon}(1,y)|^{2} \, ds dy \\ &\leqslant C \, \varepsilon^{N} \, \|w_{\varepsilon}\|_{C^{1}(\overline{\Omega})}^{2} \,, \end{split}$$

where we have used that $\int_0^{\varepsilon} \int_{\Gamma^s} r ds dy = O(\varepsilon^N)$.

The following estimates hold (see [12]), for some C > 0,

$$\|w_{\varepsilon}\|_{\mathcal{C}(\overline{\Omega})} \leqslant \frac{C}{|\lambda|+1} \|f_{\varepsilon}\|_{\mathcal{C}(\overline{\Omega})}$$
 (7.14)

$$\|w_{\varepsilon}\|_{\mathcal{C}^{1}(\overline{\Omega})} \leqslant \frac{C}{|\lambda|^{1/2} + 1} \|f_{\varepsilon}\|_{\mathcal{C}(\overline{\Omega})}. \tag{7.15}$$

Using (7.15) we have that

$$||K_{\varepsilon}||_{H^{1}(R_{\varepsilon})} \leqslant C \frac{\varepsilon^{N/2}}{|\lambda|^{1/2} + 1} ||f_{\varepsilon}||_{\mathcal{C}(\overline{\Omega})}.$$

$$(7.16)$$

which shows the first inequality of (7.8).

On the other hand we also have that

$$\left\|\lambda(A_{\varepsilon}+\lambda)^{-1}K_{\varepsilon}\right\|_{H^{1}(\Omega_{\varepsilon})} \leqslant |\lambda| \frac{1}{|\lambda|^{1/2}+1} \left\|K_{\varepsilon}\right\|_{L^{2}(R_{\varepsilon})} \leqslant C \frac{\varepsilon^{N/2}}{|\lambda|^{1/2}+1} \left\|f_{\varepsilon}\right\|_{\mathcal{C}(\overline{\Omega})}$$
(7.17)

and

$$\begin{split} \|(I - \lambda (A_{\varepsilon} + \lambda)^{-1}) (E_{\varepsilon}^{\mathcal{C}} - E_{\varepsilon}) (A_{0} + \lambda)^{-1} M f_{\varepsilon} \|_{H^{1}(\Omega) \oplus H^{1}(R_{\varepsilon})} \\ & \leq \|K_{\varepsilon}\|_{H^{1}(\Omega) \oplus H^{1}(R_{\varepsilon})} + \|\lambda (A_{\varepsilon} + \lambda)^{-1} K_{\varepsilon}\|_{H^{1}(\Omega_{\varepsilon})} \\ & \leq C \frac{\varepsilon^{N/2}}{|\lambda|^{1/2} + 1} \|f_{\varepsilon}\|_{\mathcal{C}(\overline{\Omega})} \,. \end{split}$$

Proof of Lemma 7.4: It follows easily from the definition of the extension E_{ε} and of the projection M_{ε} , that $||E_{\varepsilon}||_{\mathcal{L}(L^{\infty}(\Omega)\oplus L^{\infty}(0,1),L^{\infty}(\Omega_{\varepsilon}))}=1$ and $||M_{\varepsilon}||_{\mathcal{L}(L^{\infty}(\Omega_{\varepsilon}),L^{\infty}(\Omega)\oplus L^{\infty}(0,1))}\leqslant 1$. Hence,

$$\left\| E_{\varepsilon} A_0 (A_0 + \lambda)^{-1} M \right\|_{\mathcal{L}(L^{\infty}(\Omega_{\varepsilon}), L^{\infty}(\Omega_{\varepsilon}))} \leqslant C \left\| A_0 (A_0 + \lambda)^{-1} \right\|_{\mathcal{L}(L^{\infty}(\Omega_{\varepsilon}) \oplus L^{\infty}(0.1))}$$
(7.18)

Let $f = (f_{\Omega}, f_{R_0}) \in \mathcal{C}(\overline{\Omega}) \oplus L^{\infty}(0, 1)$, be such that

$$(A_0 + \lambda)^{-1} f = (w, v). \tag{7.19}$$

or equivalently

$$\begin{cases}
-\Delta w + \lambda w = f_{\Omega}, & \text{in } \Omega \\
\frac{\partial w}{\partial n} = 0, & \text{in } \partial \Omega \\
-\frac{1}{g} (gv_s)_s + \lambda v = f_{R_0}, & \text{in } (0, 1) \\
v(0) = w(0), \quad v(1) = w(1)
\end{cases}$$
(7.20)

proceeding as in the proof of Proposition 3.2 (iv), we have that

$$||w||_{\mathcal{C}(\overline{\Omega})} \leqslant \frac{C}{|\lambda|+1} ||f_{\Omega}||_{\mathcal{C}(\overline{\Omega})}, \quad ||v||_{\mathcal{C}(\overline{\Omega})} \leqslant \frac{C}{|\lambda|+1} \left(||f_{\Omega}||_{\mathcal{C}(\overline{\Omega})} + ||f_{R_0}||_{\mathcal{C}(0,1)} \right).$$

Since $A_0(A_0 + \lambda)^{-1} = I - \lambda(A_0 + \lambda)^{-1}$, then

$$\begin{aligned} \|A_0(A_0 + \lambda)^{-1} f\|_{\mathcal{C}(\overline{\Omega}) \oplus L^{\infty}(0,1)} &= \|f - \lambda (A_0 + \lambda)^{-1} f\|_{\mathcal{C}(\overline{\Omega}) \oplus L^{\infty}(0,1)} \\ &\leqslant \|f\|_{\mathcal{C}(\overline{\Omega}) \oplus L^{\infty}(0,1)} + C \|f\|_{\mathcal{C}(\overline{\Omega}) \oplus L^{\infty}(0,1)} \\ &\leqslant \widetilde{C} \|f\|_{\mathcal{C}(\overline{\Omega}) \oplus L^{\infty}(0,1)} \,. \end{aligned}$$

Applying this to (7.18), we have that

$$||E_{\varepsilon}A_0(A_0 + \lambda)^{-1}M||_{\mathcal{L}(L^{\infty}(\Omega_{\varepsilon}), L^{\infty}(\Omega_{\varepsilon}))} \leqslant C, \tag{7.21}$$

where C is independent of λ and ε .

Part (ii) is immediate from the fact that A_{ε} is positive and self-adjoint.

Proof of Lemma 7.5: The proof follows from Lemma 7.2, Lemma 7.3, Lemma 7.4 and statement (3.8) from Proposition 3.3.

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